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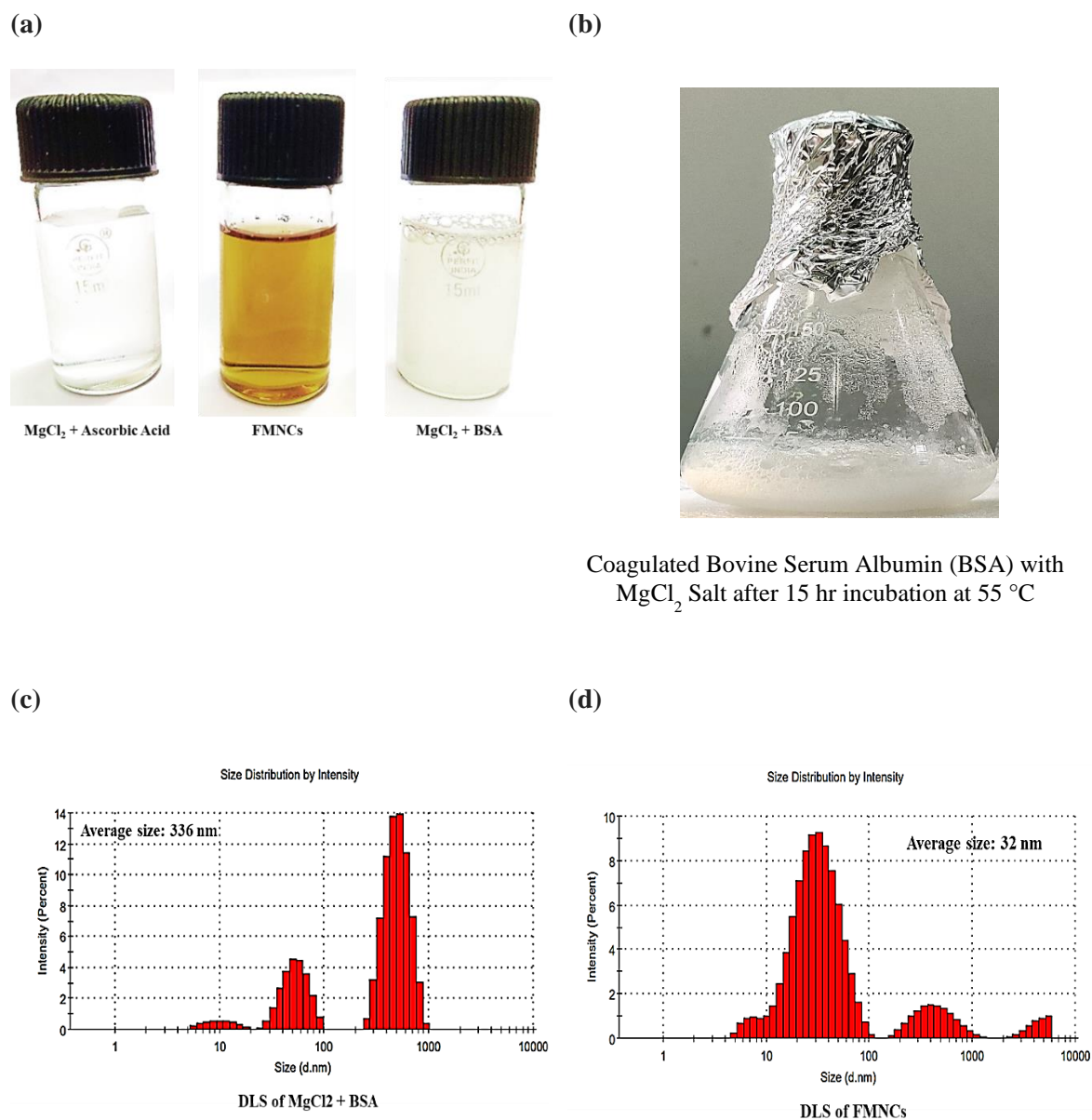
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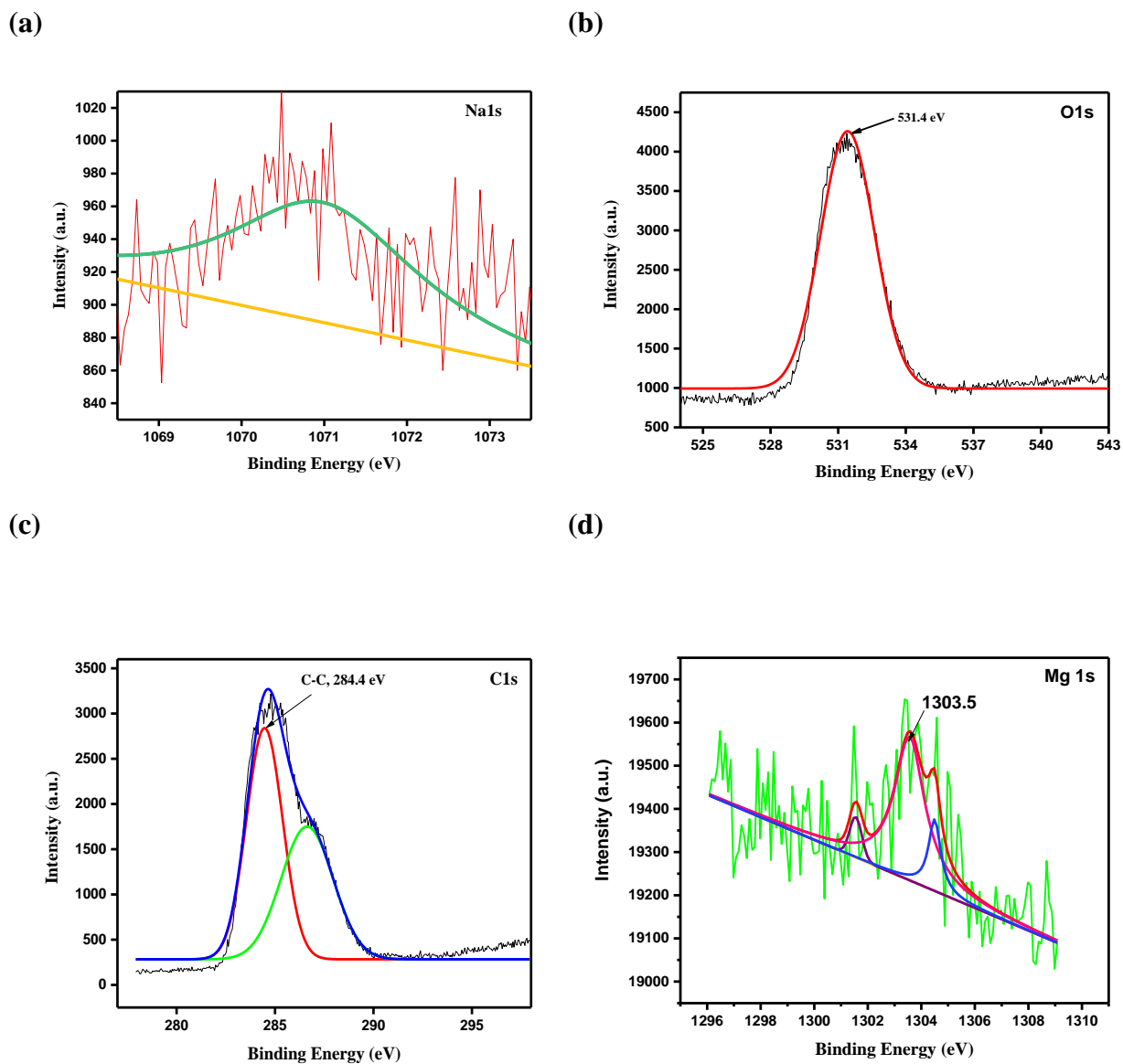
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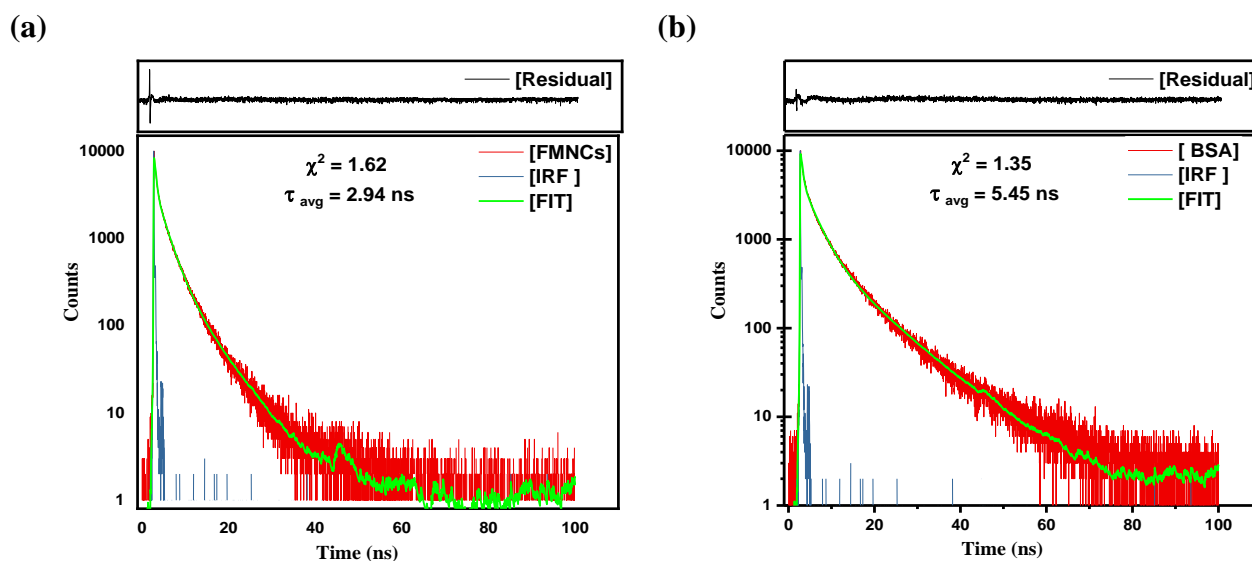
## Appendix A:



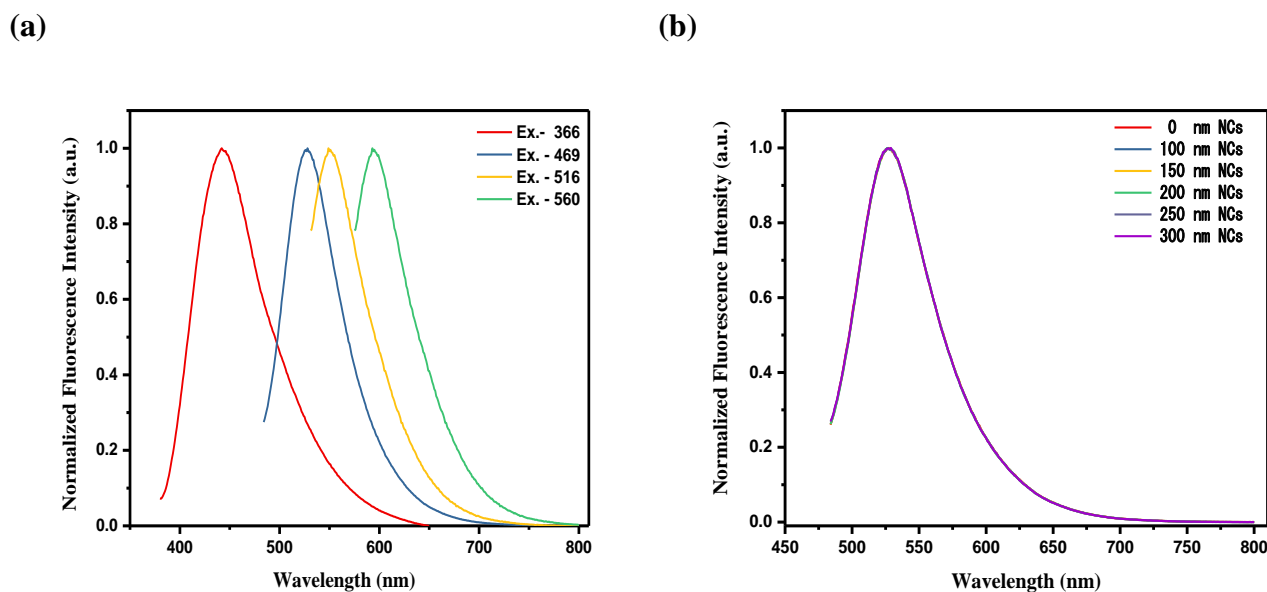
**Figure. A1** (a) Image of FMNCs and Control 1 (Mgcl<sub>2</sub>+ BSA) and control 2 (Mgcl<sub>2</sub>+Ascorbic acid), (b) Image of incubated Mgcl<sub>2</sub> + BSA at 55 °C, (C) DLS of control1 (Mgcl<sub>2</sub>+ BSA), (d) DLS of FMNCs.



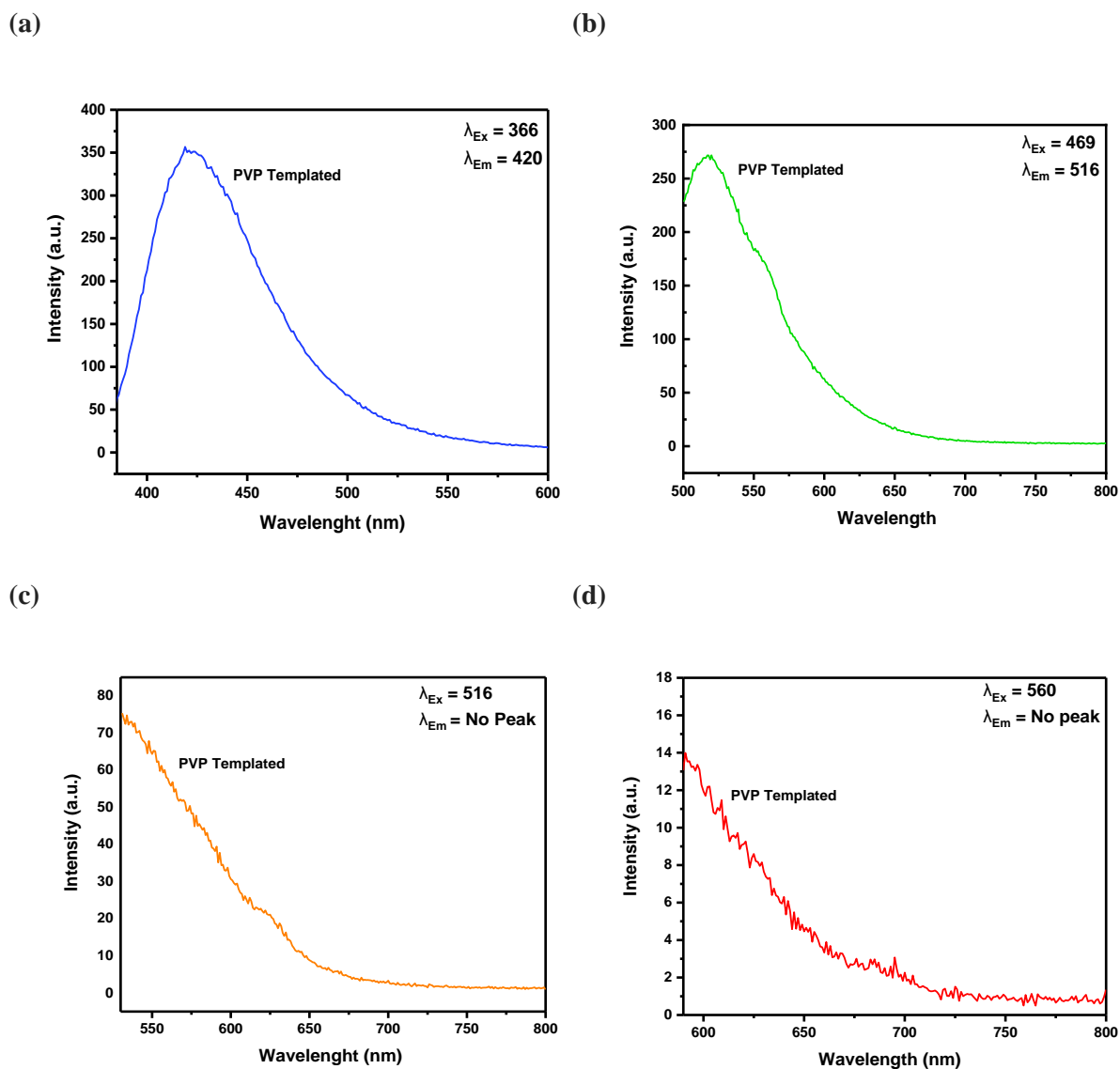
**Figure. A2** XPS spectra of (a) Na 1s, (b) O 1s, (c) C 1s, (d) Mg 1s



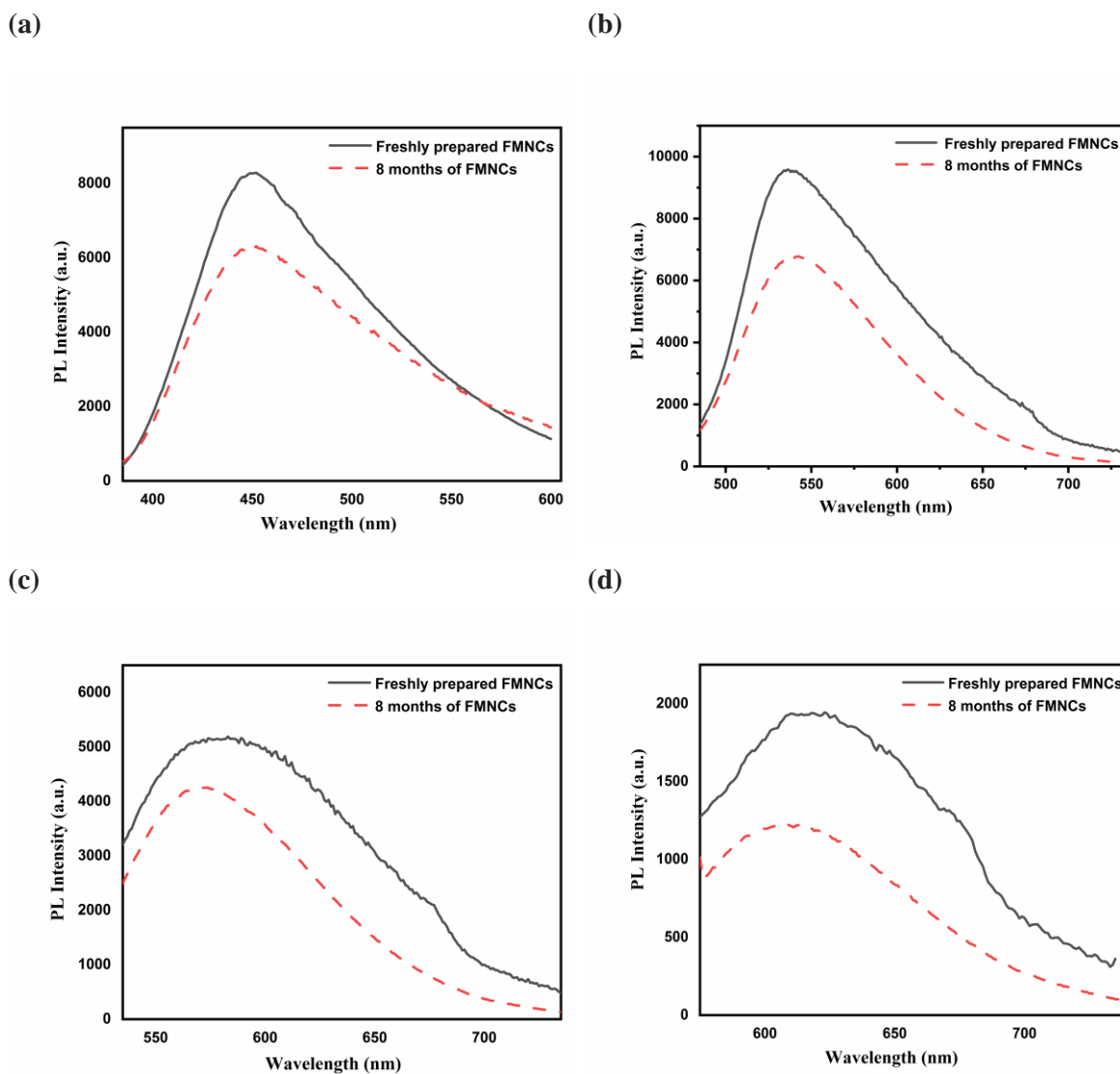
**Figure. A3** (a) Fluorescence life time decay profile of FMNCs ( $\lambda_{Ex}$  375 nm/ $\lambda_{Em}$  459 nm). (b) Fluorescence life time decay profile of processed BSA protein under the same excitation value ( $\lambda_{Ex}$  375 nm/ $\lambda_{Em}$  459 nm)



**Figure. A4** (a) Normalized Fluorescence emission spectra of lyophilized FMNCs powder (one year old) at different excitations (366,469, 516 and 560). (b) Normalized fluorescence emission spectra of FMNCs after addition of various concentration of freshly prepared NaCl solution at  $\lambda_{Ex}$  469.



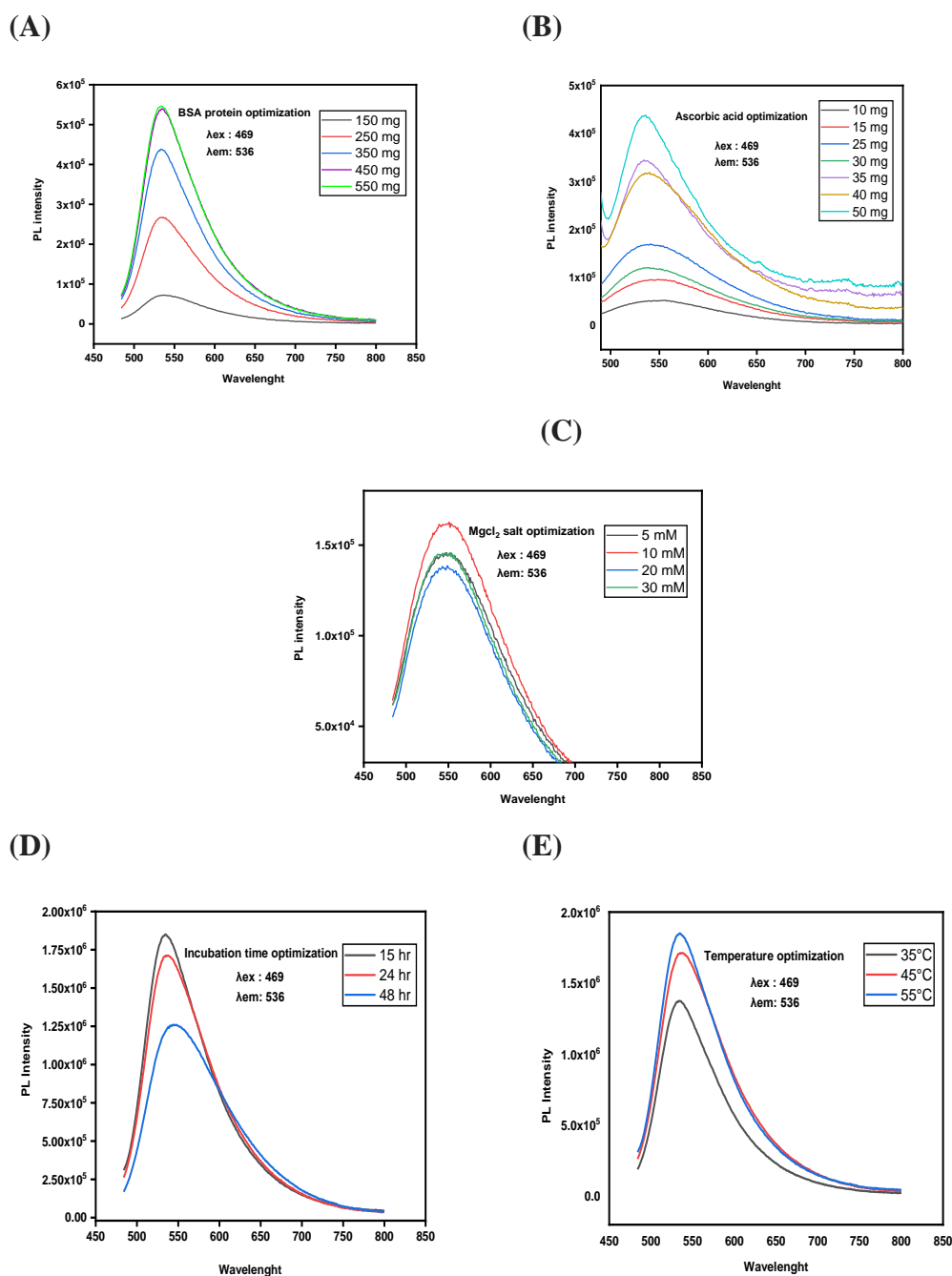
**Figure. A5** (a) Fluorescence spectra of PVP templated FMNCs nanocluster at  $\lambda_{Ex}$  366, (b) Fluorescence spectra of PVP templated FMNCs nanocluster at  $\lambda_{Ex}$  469, (c) Fluorescence spectra of PVP templated FMNCs nanocluster at  $\lambda_{Ex}$  516, (d) Fluorescence spectra of PVP templated FMNCs nanocluster at  $\lambda_{Ex}$  560.



**Figure. A6** (a) Fluorescence emission spectra of 8 months old and freshly prepared aqueous FMNCs at  $\lambda_{Ex}$  366 (b) Fluorescence emission spectra of 8 months old and freshly prepared aqueous FMNCs at  $\lambda_{Ex}$  469 (c) Fluorescence emission spectra of 8 months old and freshly prepared aqueous FMNCs at  $\lambda_{Ex}$  516 (d) Fluorescence emission spectra of 8 months old and freshly prepared aqueous FMNCs at  $\lambda_{Ex}$  560.

## Appendix B

Figure B1



**Figure B1** (A) BSA protein optimization (450 mg/ml optimized) (B) Ascorbic acid optimization (35 mg/ml optimized) (C)  $MgCl_2$  salt optimization (10mM optimized) (D) Incubation time optimization (15 hr. optimized) (E) Temperature optimization (55 ° C optimized).

<u>Nucleation models</u>	<u>Model</u>	<u>Differential form f(<math>\alpha</math>)</u>	<u>Integral form g(<math>\alpha</math>)</u>
<u>P2</u>	<u>Power Law</u>	$(2/3) \alpha^{-1/2}$	$\alpha^{3/2}$
<u>P3</u>	<u>Power Law</u>	$2\alpha^{1/2}$	$\alpha^{1/2}$
<u>P4</u>	<u>Power Law</u>	$3\alpha^{2/3}$	$\alpha^{1/3}$
<u>P5</u>	<u>Power Law</u>	$4\alpha^{3/4}$	$\alpha^{1/4}$
<b><u>Sigmoidal rate equations</u></b>			
<u>A1</u>	<u>Avrami-Erofeev</u>	$(3/2) (1-\alpha) [-\ln(1-\alpha)]^{1/3}$	$[-\ln(1-\alpha)]^{2/3}$
<u>A2</u>	<u>Avrami-Erofeev</u>	$2(1-\alpha) [-\ln(1-\alpha)]^{1/2}$	$[-\ln(1-\alpha)]^{1/2}$
<u>A3</u>	<u>Avrami-Erofeev</u>	$3(1-\alpha) [-\ln(1-\alpha)]^{2/3}$	$[-\ln(1-\alpha)]^{1/3}$
<u>A4</u>	<u>Avrami-Erofeev</u>	$4(1-\alpha) [-\ln(1-\alpha)]^{3/4}$	$[-\ln(1-\alpha)]^{1/4}$
<u>F1</u>	<u>Prout-Tompkins</u>	$\alpha(1-\alpha)$	$\ln[\alpha/(1-\alpha)]$
<u>F2</u>	<u>Contracting area</u>	$2(1-\alpha)^{1/2}$	$1-(1-\alpha)^{1/2}$
<u>F3</u>	<u>Contracting volume</u>	$3(1-\alpha)^{2/3}$	$1-(1-\alpha)^{1/3}$
<u>F4</u>	<u>Random nucleation (1)</u>	$(1-\alpha)^2$	$1/(1-\alpha)$
<u>F5</u>	<u>Random nucleation (2)</u>	$(1-\alpha)^3/2$	$1/(1-\alpha)^2$
<b><u>Diffusion models</u></b>			
<u>D1</u>	<u>1D diffusion</u>	$1/2\alpha$	$\alpha^2$
<u>D2</u>	<u>2D diffusion-Valensi</u>	$[-\ln(1-\alpha)]^{-1}$	$(1-\alpha) \ln(1-\alpha) + \alpha$
<u>D3</u>	<u>3D diffusion-jander</u>	$(3/2) (1-\alpha)^{2/3} / [1-(1-\alpha)^{1/3}]$	$[1-(1-\alpha)^{1/3}]^2$
<u>D4</u>	<u>3D diffusion-Ginstling</u>	$(3/2)/[(1-\alpha)^{-1/3} - 1]$	$1-2\alpha/3-(1-\alpha)^{2/3}$
<b><u>Reaction order models</u></b>			
<u>R1</u>	<u>First order</u>	$1-\alpha$	$-\ln(1-\alpha)$
<u>R2</u>	<u>Second order</u>	$(1-\alpha)^2$	$(1-\alpha)^{-1} - 1$
<u>R3</u>	<u>Third order</u>	$(1-\alpha)^3$	$[(1-\alpha)^{-2} - 1]/2$
<u>R4</u>	<u>One and half order</u>	$(1-\alpha)^{3/2}$	$2[(1-\alpha)^{-1/2} - 1]$

Table B1

**Table B2** Preexponential kinetic factor (A) at 10 °C/min, computed by utilizing KAS-derived equation.

Conversion( $\alpha$ )	A ( $\text{min}^{-1}$ )	A (Nuclei $\mu\text{m}^{-2} \text{min}^{-1}$ )
0.1	1.89077E+14	1.25645E+30
0.2	1.80698E+13	1.20077E+29
0.3	4.35197E+17	2.89E+33
0.4	3.61E+16	2.4E+32
0.5	7.8E+14	5.18E+30
0.6	1.29E+13	8.58E+28
0.7	9.63E+09	6.4E+25
0.8	4.92E+15	3.27E+31
0.9	778517597.8	5.1734E+24

**Table B3** Nucleation rate computed by utilizing kinetic and thermodynamic barriers, essential for random nucleation process within MgNCs matrix.

T(K)	<i>J</i> (nuclei $\mu\text{m}^{-2} \text{min}^{-1}$ ) 10°C/min	T(K)	<i>J</i> (nuclei $\mu\text{m}^{-2} \text{min}^{-1}$ ) 15°C/min	T(K)	<i>J</i> (nuclei $\mu\text{m}^{-2} \text{min}^{-1}$ ) 20°C/min
533.15	0.158867	539.22	0.4431	543.29	0.6007
570.19	16.63	576.16	42.191	579.24	45.84
587.18	128.59	594.24	351.78	597.24	376.76
600.17	451.47	608.17	1298.47	612.24	1570.89
616.23	1639.61	626.26	5360.72	630.29	6459.00
644.15	7461	658.24	29338.92	664.24	41301.41
697.21	5221919	705.16	11849271	712.16	16233076
777.18	19799076	795.16	65536129	807.21	1.11×10 <sup>8</sup>

**Table B4** Value of exponential function constant  $Q_1$ ,  $Q_2$ ,  $\rho_1$ , and  $\rho_2$ , acquired from a theoretical model by fitting the experimental data.

Exponential function constants	10 °C/min	15 °C/min	20 °C/min
$Q_1$	587.11	918.573	879.90
$\rho_1$	58.37	51.31	48.69
$Y_0$	-898778.80	-1648835.36	-1985815.90
$Q_2$	28.11	18.27	11.12
$\rho_2$	0.071±0.00642	0.05749±0.0029	0.05169±0.00191
$J_0$	-191008.20	-315372.41	-358593.10

\*\*  $Q_1$  and  $\rho_1$  are exponential functions constant and  $Y_0$  is intercept for a theoretical model of nucleation vs temperature, whereas  $Q_2$ ,  $\rho_2$  and are exponential function constant and  $J_0$  is intercept for nucleation vs conversion.

**Table B 5** Interfacial energy was calculated by using a thermodynamic barrier within the MgNCs matrix with respect to temperature.

T(K)	Y (mJ/m <sup>2</sup> )- 10 <sup>0</sup> C/min	T(K)	Y (mJ/m <sup>2</sup> )- 15 <sup>0</sup> C/min	T(K)	Y (mJ/m <sup>2</sup> )- 20 <sup>0</sup> C/min
473.21	111.75	479.22	112.81	481.25	113.03
533.15	114.66	539.22	115.388	543.29	116.09
570.19	119.56	576.16	120.23	579.24	120.79
587.18	122.00	594.24	122.82	597.24	123.36
600.17	123.93	608.17	124.87	612.24	125.56
616.23	126.30	626.26	127.50	630.29	128.19
644.15	130.44	658.24	132.18	664.24	133.13
697.21	136.88	705.16	137.74	712.16	138.80
777.18	148.01	795.16	150.09	807.21	151.78

### Section B1 Calculation of $\alpha$ – helix content using CD Spectra:

The  $\alpha$ -helix content of BSA and FMNCs is calculated by using the following equation: -

$$\% \text{ helix} = \left( \frac{\theta_{mrd} - 4000}{33000 - 4000} \right) \times 100$$

$\theta_{mrd}$  is the mean molar ellipticity per residue (deg cm<sup>2</sup>/dmol)

$$\theta_{mrd} = \frac{\theta_d M}{10 ncl}$$

$\theta_d$  is ellipticity in mdeg, M is the molecular weight of BSA in Da, c is concentration in mg/ml and l is path length in cm.

In this particular calculation:

M= 66050 Da, c= 0.5mg/ml (Pure BSA) and 0.4 mg/ml (MgNCs) and n is 583 (Amino acid residue in BSA)

The calculated  $\alpha$  helicity was found to be 57.7% for Pure BSA, which becomes 22% after MgNCs formation, this result shows a drastic change in the conformation of BSA after MgNCs formation.

**Section B2** Calculation  $\Delta G^*$  using CNT (Classical nucleation Theory) and Thermogravimetric analysis (TGA):

The Nucleation Rate for thermally activated process can be expressed as (*J. W. Mullin, Crystallization, Butterworth-Heinemann, Oxford, 1993*)(Mullin):

$$J = A \exp(\Delta G^*/kT) \quad (1)$$

Where,  $k$  is Boltzmann constant ( $1.3805 \times 10^{-23} \text{JK}^{-1}$ ),  $T$  is temperature, and  $\Delta G^*$  is critical free energy barrier to nucleation.

We assume that,

$$\text{Per molecule gas constant in 1 Mg cluster} = R/N_A = k \quad (2)$$

Where,  $R$  is universal gas constant ( $8.314 \text{ J K}^{-1} \text{ mol}^{-1}$ ), and  $N$  is Avogadro number ( $6.023 \times 10^{23} \text{ mol}^{-1}$ ). (*J. W. Mullin, Crystallization, Butterworth-Heinemann, Oxford, 1993*)(Mullin):

Hence,

$$\text{We can write} \quad N_A = R \quad (3)$$

$$6.023 \times 10^{23} \text{ mol}^{-1} = R \quad (4)$$

Now,  $6.023 \times 10^{23} \text{ mol}^{-1}$  Mg clusters are nucleating at particular temperature ( $T$ ) and respective conversion ( $\alpha$ ). So from eq. 2 we can write

$$6.023 \times 10^{23} \text{ mol}^{-1} \rightarrow \frac{R}{N_A} \times N_A = k \quad (5)$$

Hence,

$$1 \text{ Mole of Mg Clusters} \rightarrow k = R \quad (6)$$

Putting the value of eq (6) in eq. (1) for nucleation of 1 mole of the particles eq. (7) is derived below.

$$J = A \exp(\Delta G^*/RT) \quad (7)$$

But, in eq. (1) and eq. (7), while computation of nucleation rate, only  $A$  ( Pre-exponential factor) is used during computation of kinetic barrier of nucleation, and activation energy of nucleation ( $E_\alpha$ ) is still missing. Thus, to compute nucleation rate, both kinetic and thermodynamic barriers were taken into account and expressed as below(Hu et al., 2012):

$$J = A \exp\left(\frac{-E}{kT}\right) \exp\left(\frac{-\Delta G^*}{kT}\right) \quad (8)$$

To convert eq. (8) in terms of R, we will have to use  $k=R$  from eq. (6) thus by using eq. (6) in eq. (8), the nucleation rate equation gets converted in terms of R as follows:’

$$J = A \exp\left(\frac{-E}{RT}\right) \exp\left(\frac{-\Delta G^*}{RT}\right) \quad (9)$$

Same eq. (9) is used to calculate nucleation rate by Qingyun Li et.al.

Thus, in current research, we are using eq. (9) for computation of nucleation rate. As our research goal is to calculate nucleation rate by TGA iso-conversional models at high temperature, the  $E$  is converted to  $E_\alpha$ ,  $A$  is converted to  $A_\alpha$  and  $\Delta G^*$  is also computed iso-conversionally, and finally eq. (9) can be expressed as by equation below:

$$J = A_\alpha \exp\left(\frac{E_\alpha}{RT}\right) \exp\left(\frac{\Delta G^*}{RT}\right) \quad (10)$$

The thermodynamic barrier  $\Delta G^*$  can be expressed as below: (J. W. Mullin, Crystallization, Butterworth-Heinemann, Oxford, 1993)

$$\Delta G^* = \frac{16\pi\gamma^3 v^2}{3(kT \ln S)^2} \quad (11)$$

After putting eq. (11) into eq. (8), we get,  $J$  as

$$J = A \exp\left(\frac{-E}{kT}\right) \exp\left(\frac{-16\pi\gamma^3 v^2}{3k^3 T^3 (\ln S)^2}\right) \quad (12)$$

In eq. (12), the thermodynamic barrier to nucleation rate is expressed as

$$B^* = \frac{16\pi v^2 \gamma^3}{3k^3 T^3 (\ln S)^2} \quad (13)$$

In our study, by examining  $z(\alpha)$  master plots, it is already verified that our nucleation is random. In random nucleation process, during the formation of first nuclei, the nucleation is observed only, when the nuclei size ( $r$ ) is equal to the standard critical size of nucleation ( $r^*$ ), according to the equation below (Dirksen and Ring, 1991):

$$r^* = \frac{2\beta_a \gamma v}{3\beta_v k_B T \ln(S)} \quad (14)$$

Where,  $\beta_a$  is  $4\pi$ ,  $\beta_v$  is  $4\pi/3$ ,  $\gamma$  is surface free energy per unit area,  $v$  is volume of aggregate,  $S$  is supersaturation, and  $T$  is temperature. When radius of nuclei formed becomes equal to the standard critical size of nucleation, as per eq. (13), the first nuclei formation takes place at particular temperature and the value of supersaturation ( $S$ ) becomes 2.718 (Dirksen and Ring, 1991).

In our research, we are observing nuclei formation by R3 mechanism ( $z(\alpha)$  master plots), as a starting phase of nucleation (when actual nucleation starts), thus in our case the value of supersaturation,  $S$  is 2.718.

By putting the value of  $S=2.718$ , the thermodynamic term  $B^*$  in eq. (13) becomes,

$$B^* = \frac{16\pi v^2 \gamma^3}{3k^3 T^3 (\ln 2.718)^2}$$

$$B^* = \frac{16\pi v^2 \gamma^3}{3k^3 T^3 (1)^2}$$

$$B^* = \frac{16\pi v^2 \gamma^3}{3k^3 T^3} \quad (15)$$

This eq. (15) is also used as a thermodynamic term, during estimation of nucleation kinetics of a L-glutamic acid (Lindenberg and Mazzotti, 2009). This thermodynamic term in eq. (15) is used to estimate interfacial energy variation with temperature.

The unit of R.H.S. of eq. (15) is

$$\frac{16\pi v^2 \gamma^3}{3k^3 T^3} = \frac{(m^3)^2 \times \left(\frac{J}{m^2}\right)^3}{\left(\frac{J}{K}\right)^3 \times K^3} = (\text{unitless})$$

Thus, we adopted L.H.S of eq. (15) to be as.

$$\frac{\Delta G^*}{RT} = \frac{\left(\frac{J}{mol}\right)}{\left(\frac{J}{mol K}\right) \times K} = (\text{unitless})$$

Thus, eq. (15) is expressed as,

$$\frac{\Delta G^*}{RT} = \frac{16\pi v^2 \gamma^3}{3k^3 T^3} \quad (16)$$

Thus, the eq. (16) is used to calculate interfacial energy against temperature.

The current nucleation of ultrasmall magnesium clusters is highly dependent on the temperature, and thereby TGA technology is used to understand the nucleation kinetics of ultrasmall magnesium clusters. During calculation of thermodynamic barrier ( $\Delta G^*$ ), eq. (11) is used. The calculated value of ( $\Delta G^*$ ) from eq. (11) is called as a free energy barrier to the nucleation at high temperature (Singh et al., 2021b). As this  $J$  in eq. (10) is used to compute nucleation rate against conversion (0.1 to 0.9), thus  $E_\alpha$  and  $\Delta G^*$  need to be calculated against conversion varying from 0.1 to 0.9 by TGA. Thus, by TGA for the calculation of free energy barrier to nucleation ( $\Delta G^*$ ) against conversion is used. To compute interfacial energy, eq. (16) is used:

$$\begin{aligned} \frac{\Delta G^*}{RT} &= \frac{16\pi v^2 \gamma^3}{3k^3 T^3} \\ \frac{3k^3 T^3 \Delta G^*}{16RT\pi v^2} &= \gamma^3 \\ \gamma &= \left(\frac{3k^3 T^3 \Delta G^*}{16RT\pi v^2}\right)^{1/3} \quad (17) \end{aligned}$$

At 0.5 conversion and corresponding temperature of 600.17 K, eq. (17) can be computed as follows,

$$\gamma = \left(\frac{3 \times (1.38 \times 10^{-20})^3 \times (600.17)^3 \times 1.51 \times 10^8}{16 \times 8314 \times 600.17 \times 3.14 \times (2.32 \times 10^{-29})^2}\right)^{1/3}$$

$$\begin{aligned}
&= \left( \frac{2.5737 \times 10^{-43}}{1.3493 \times 10^{-49}} \right)^{1/3} \\
&= (1907433.484)^{1/3} \\
&= 124.017 \text{ mJ/m}^2
\end{aligned}$$

At 600.17 K, from eq. (10), according to Classical Nucleation Theory (CNT), we get

$$\begin{aligned}
(B^*)_{CNT} &= \frac{16\pi v^2 \gamma^3}{3k^3 T^3 (\ln S)^2} \\
\frac{\Delta G^*_{CNT}}{RT} &= \frac{16\pi v^2 \gamma^3}{3k^3 T^3 (\ln S)^2} \\
\Delta G^*_{CNT} &= \frac{16\pi v^2 \gamma^3 RT}{3k^3 T^3 (\ln S)^2}
\end{aligned}$$

At 600.17 K,  $\Delta G^*_{CNT}$  is expressed as,

$$\begin{aligned}
\Delta G^*_{CNT} &= \frac{16\pi v^2 \gamma^3 RT}{3k^3 T^3 (\ln S)^2} \\
&= \frac{16 \times 3.14 \times (2.32 \times 10^{-29})^2 \times (124.017)^3 \times 8314 \times 600.17}{3 \times (1.38 \times 10^{-20})^3 \times (600.17)^3 \times (\ln 2.718)^2} \\
&= \frac{2.573674 \times 10^{-43}}{1.704438 \times 10^{-51}} \\
&= 150998393.61 \text{ mJ/mol} \\
\Delta G^*_{CNT} &= 150.998 \text{ kJ/mol}
\end{aligned}$$

At 600.17 K, Free energy change of nucleation by TGA ( $\Delta G^*_{TGA}$ ) is expressed as,

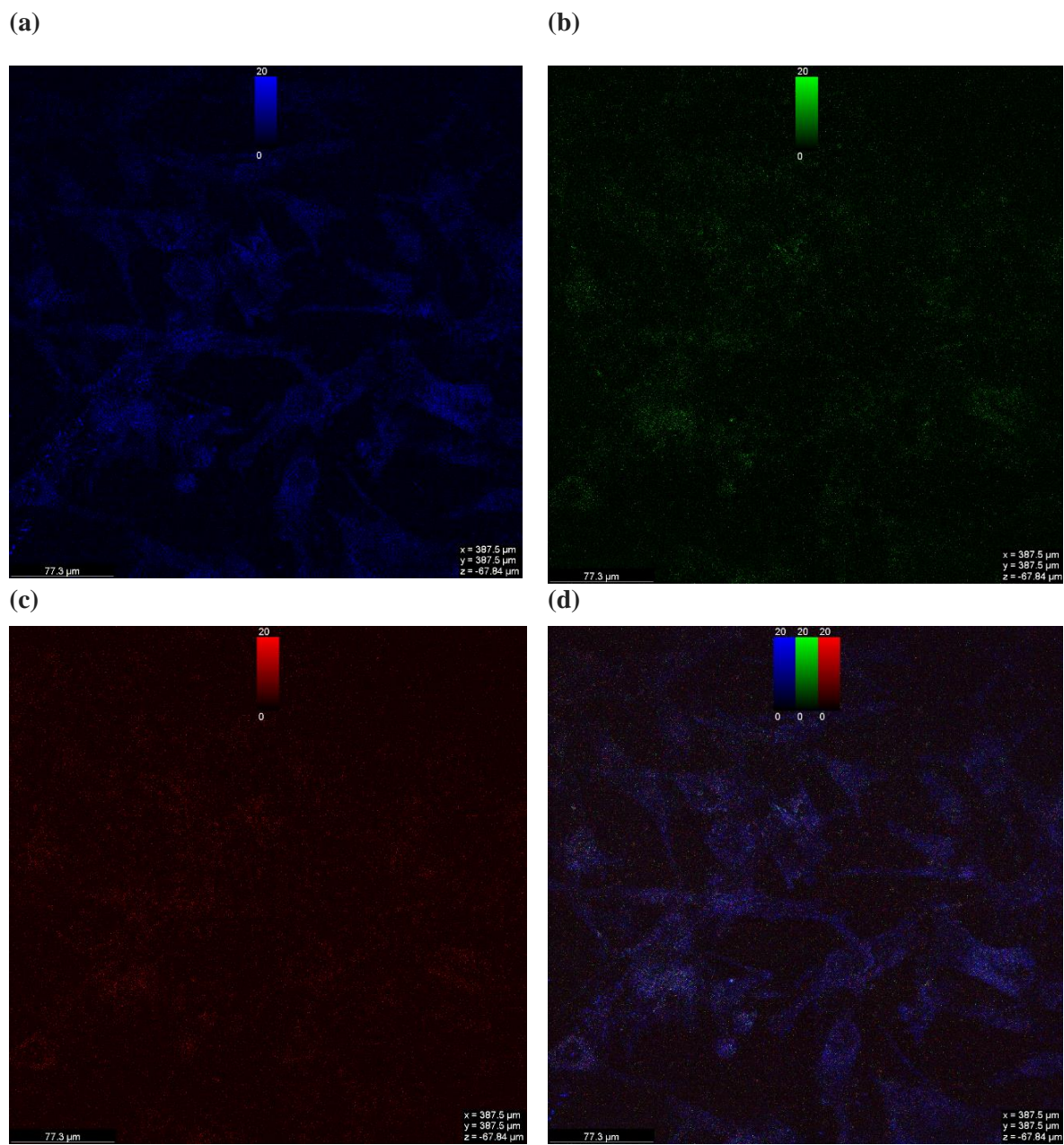
$$\begin{aligned}
\Delta G^*_{TGA} &= E_\alpha + RT_p \ln \left( \frac{K_B T_p}{hA} \right) \\
&= E_\alpha + RT_p \ln \left( \frac{K_B T_p}{hA} \right)
\end{aligned}$$

Where  $E_\alpha$  = Activation energy at particular conversion in Joules/mole,  $k$  is Boltzmann constant ( $1.3805 \times 10^{-23} \text{ JK}^{-1}$ ),  $R$  molar gas constant ( $8.314 \text{ Joules/K-mole}$ )

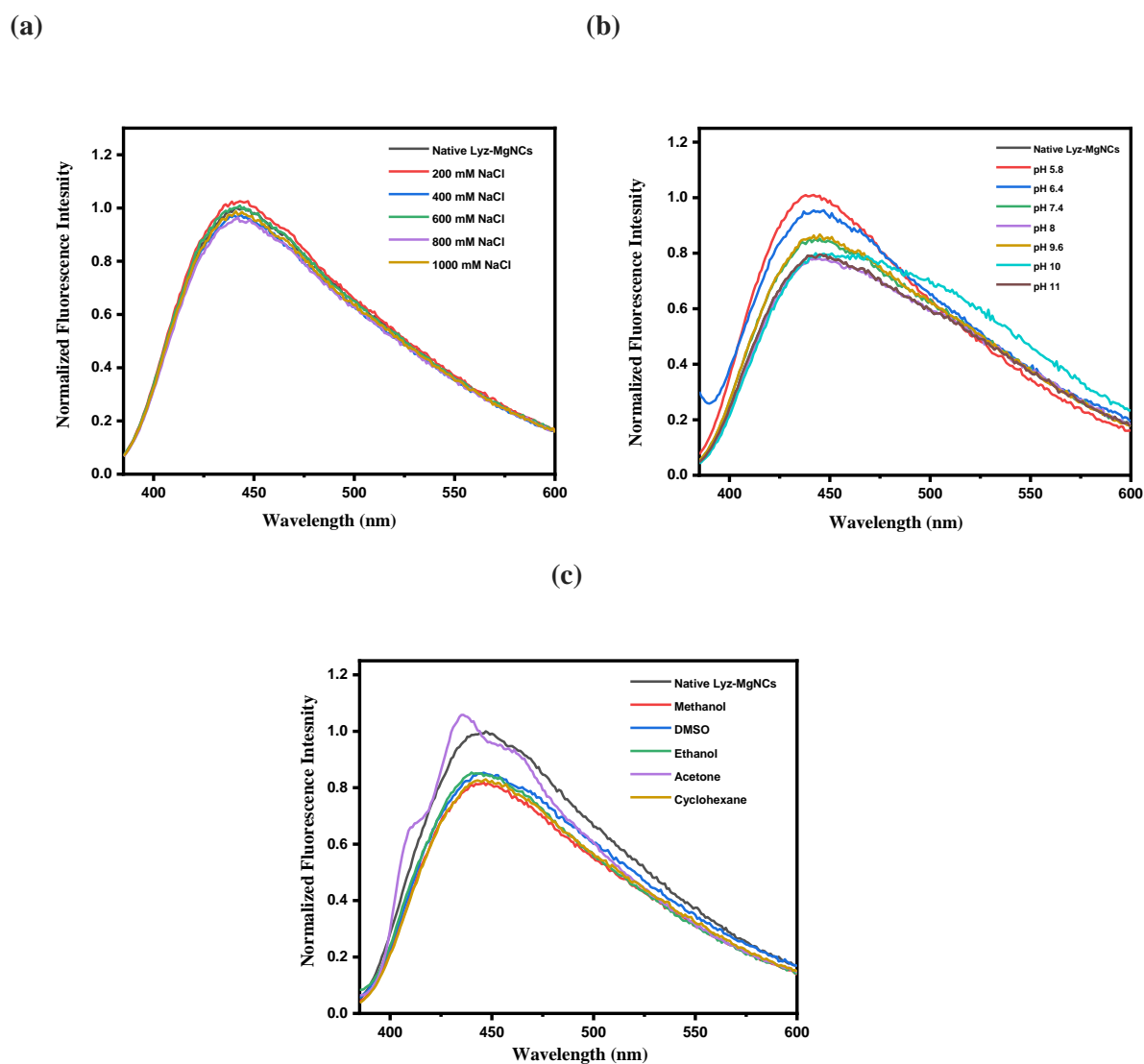
$$\begin{aligned}
&= 171401 + 8.314 \times 592.05 \ln \left( \frac{1.380649 \times 10^{-23} \times 592.05}{(6.626 \times 10^{-34}) \times (7.8 \times 10^{14})} \right) \\
&= 171401 + 8.314 \times 592.05 \ln \left( \frac{8.17413 \times 10^{-21}}{5.16828 \times 10^{-19}} \right) \\
&= 171401 + 8.314 \times 592.05 \times \ln 0.0158159 \\
&= 171401 + 8.314 \times 592.05 \times (-4.1467) \\
&= 171401 - 20411.316 \\
&= 150989 \text{ J/mol} \\
\Delta G^*_{TGA} &= 150.989 \text{ kJ/mol}
\end{aligned}$$

Hence it is proved at 0.5 conversion and 600.17K,  $\Delta G^*$  calculated from Classical Nucleation Theory ( $\Delta G^*_{CNT} = 150.998 \text{ kJ/mol}$ ) and TGA technique ( $\Delta G^*_{TGA} = 150.989 \text{ kJ/mol}$ ) comes out to be same.

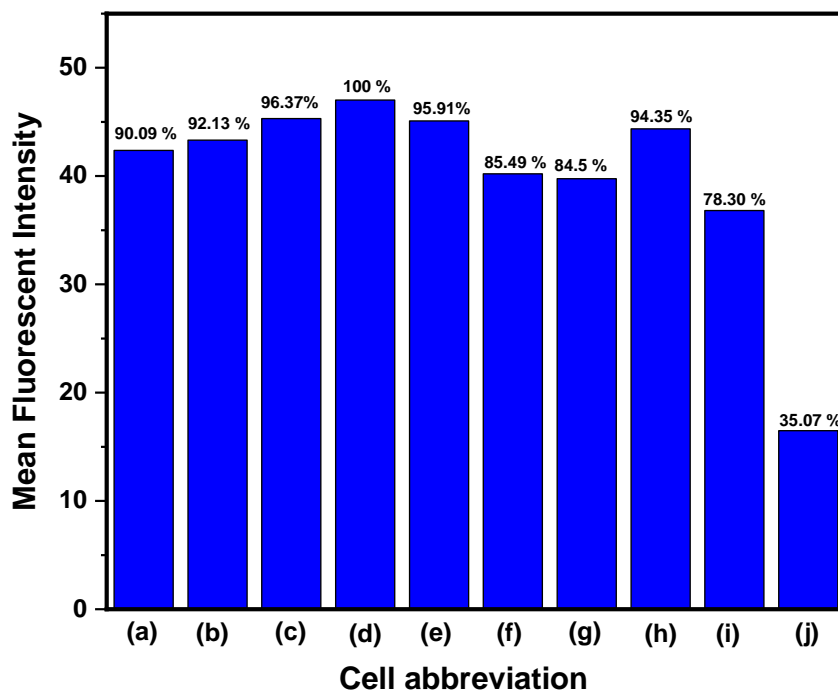
## Appendix C



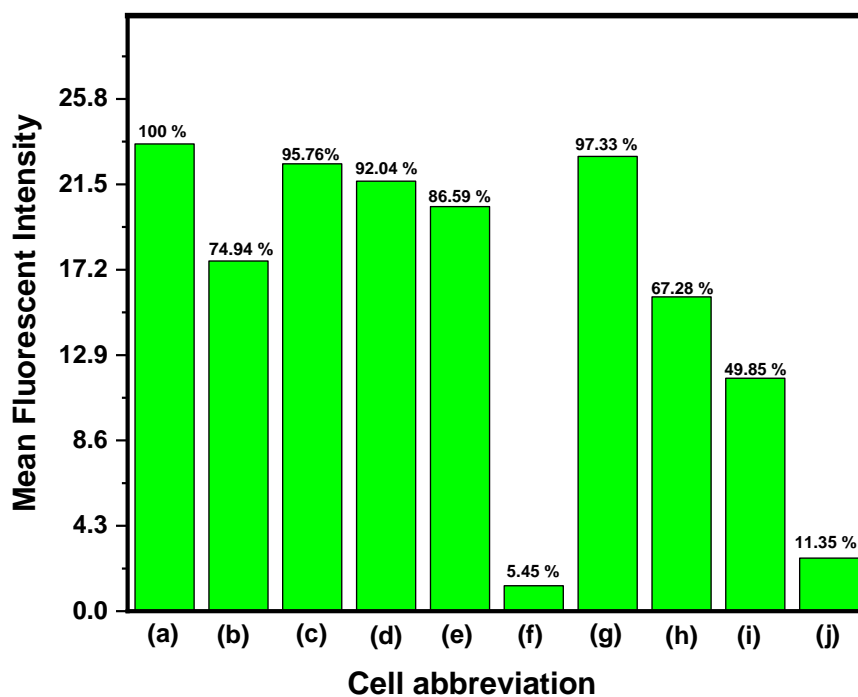
**Figure. C1** Multifluorescent confocal image of U 87-MG cells treated with Lysozyme.



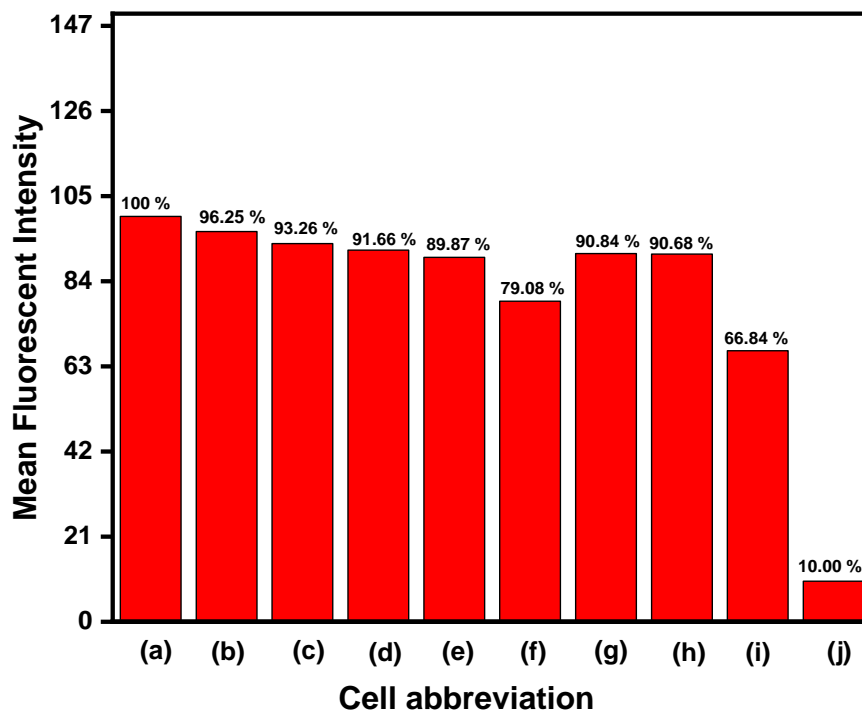
**Figure. C2** (a) Normalized fluorescence emission spectra of Lyz-MgNCs after addition of various concentration of freshly prepared NaCl solution at  $\lambda_{\text{Ex}}$  366 nm, (b) pH. Dependent Fluorescence emission spectra at different pH. (5.8-11) Values at  $\lambda_{\text{Ex}}$  366 nm, (c) Normalized fluorescence intensity of Lyz-MgNCs with different solvents at  $\lambda_{\text{Ex}}$  366 nm



**Figure. C3** Mean Luminescence Intensity in U -87 MG cells (1) Lyz-MgNCs [Ex. 366 nm, Em. Blue (450 nm)] ((a) = 0  $\mu\text{m}$ , (b) = 2 $\mu\text{m}$ , (c) = 4 $\mu\text{m}$ , (d) = 6 $\mu\text{m}$ , (e) = 8 $\mu\text{m}$ , (f) = 10 $\mu\text{m}$ , (g) = 12 $\mu\text{m}$ , (h) = 14  $\mu\text{m}$ , (i) = 16  $\mu\text{m}$ , (j) = 18  $\mu\text{m}$ ).



**Figure. C4** Mean Luminescence Intensity in U -87 MG cells Lyz-MgNCs [Ex. 469 nm, Em. Green (545 nm)] ((a) = 0  $\mu\text{m}$ , (b) = 2 $\mu\text{m}$ , (c) = 4 $\mu\text{m}$ , (d) = 6 $\mu\text{m}$ , (e) = 8 $\mu\text{m}$ , (f) = 10 $\mu\text{m}$ , (g) = 12 $\mu\text{m}$ , (h) = 14  $\mu\text{m}$ , (i) = 16  $\mu\text{m}$ , (j) = 18  $\mu\text{m}$ )



**Figure. C5** Mean Luminescence Intensity in U -87 MG cells Lyz-MgNCs [Ex. 560 nm, Em. Red ( 628 nm)] ((a) = 0  $\mu\text{m}$ , (b) = 2 $\mu\text{m}$ , (c) = 4 $\mu\text{m}$ , (d) = 6 $\mu\text{m}$ , (e) = 8 $\mu\text{m}$ , (f) = 10 $\mu\text{m}$ , (g) = 12 $\mu\text{m}$ , (h) = 14  $\mu\text{m}$ , (i) = 16  $\mu\text{m}$ , (j) = 18  $\mu\text{m}$ )

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