

Chapter 4

Generalized Sobolev type spaces involving the Weinstein transform

4.1 Introduction

Sobolev spaces are a fundamental concept in functional analysis and the theory of partial differential equations. They provide a framework for studying the properties of functions that have derivatives that may not necessarily be defined in the classical sense. The norm in a Sobolev space is defined in terms of derivatives and integrals, and it measures the smoothness of a function. It is widely used for solving the problems of partial differential equations, as well as in nonlinear analysis, differential geometry, physics, and other areas of mathematics. In [40], Park and Kang extended the notion of Sobolev space to generalized distribution spaces of Beurling-Björck type [6, 8], and used the weight function ω to investigate the Sobolev imbedding theorem, the Rellich's compactness theorem, and others. Pathak and Pandey [44] introduced the Sobolev space of type $G_{\mu}^{p,s}$ and discussed some properties, such as

completeness and inclusion, by using the definitions of the distributional Hankel transform and they introduced a new concept called the Hankel potential \mathcal{H}_μ^s and found that Hankel potential is a continuous linear mapping from the Zemanian space \mathcal{H}_μ^s into the itself. Later on, Pathak and Shrestha [45], defined the space of type $G_{\omega,\mu}^{p,s}$ and discussed many results. Pathak [43], considered the generalized Sobolev space $H_\omega^s(\mathbb{R}^n)$ which is a generalization of the Sobolev space $H^s(\mathbb{R}^n)$ and developed a multiresolution analysis for the generalised Sobolev space.

The Weinstein transform has a rich calculus and involves many problems in the mathematical sciences and their applications. Utilizing this theory, many researchers obtained observations in the theory of Sobolev-type spaces. In this connection, Salem and Dachraoui [4], studied the Sobolev type spaces $G_{\alpha,\beta}^{p,s}$ associated with Jacobi differential operators $\Delta_{\alpha,\beta}$, and examined that the Jacobi potential $\mathcal{H}_{\alpha,\beta}^s$, is a pseudodifferential operator with a precise symbol. Using the aforesaid results they extended the operator $\mathcal{H}_{\alpha,\beta}^s$, to the distributional space [8]. Hassen and Belgacem [35], introduced the Sobolev spaces of exponential type associated with the Weinstein operator, and found many results. Mehrez, [30], Mejjaoli, et al. [31–33], Salem [56], Saoudi [60, 61], and Trimèche, [81], observed useful results regarding the Weinstein transform.

In this present chapter the author made a proper framework for the generalised Sobolev type space with the help of definitions and properties of the Weinstein transform. In this chapter, the generalized Sobolev space of type $G_\omega^{p,s}$ is introduced, and with the help of the Weinstein transform, and various properties are obtained. The L^p -space of Hankel potential, $W_\omega^{m,p}(\mathbb{R}_+^{n+1})$ is considered, and many results related to the Hankel potential \mathcal{H}^k are studied by exploiting the theory of Weinstein transform.

4.2 Generalized Sobolev type space $G_\omega^{p,s}$

In this section, the generalized Sobolev space $G_\omega^{p,s}$ is defined and its various properties including completeness and inclusion are obtained by exploiting the theory of the Weinstein transform.

Theorem 4.2.1. *The Weinstein transform $\mathcal{F}_w f$ is an automorphism on $S_\omega(\mathbb{R}_+^{n+1})$.*

Proof: Let $\phi \in S_\omega(\mathbb{R}_+^{n+1})$, then first we show that $\mathcal{F}_w(\phi) \in S_\omega(\mathbb{R}_+^{n+1})$. From (3.2.21), we have

$$\begin{aligned} p_{\alpha,\lambda}[\mathcal{F}_w\phi(\xi)] &= \sup_{\xi \in \mathbb{R}_+^{n+1}} e^{\lambda\omega(\xi)} | (\Delta_{W,\beta}^n)_\xi^\alpha (\mathcal{F}_w\phi)(\xi) | \\ &= \sup_{\xi \in \mathbb{R}_+^{n+1}} e^{\lambda\omega(\xi)} | (\Delta_{W,\beta}^n)_\xi^\alpha \int_{\mathbb{R}_+^{n+1}} e^{-i\langle \xi', y' \rangle} J_\beta(\xi_{n+1} y_{n+1}) \phi(y) d\mu_\beta(y) | . \end{aligned}$$

Using (1.4.5), we get

$$\begin{aligned} p_{\alpha,\lambda}[\mathcal{F}_w\phi(\xi)] &= \sup_{\xi \in \mathbb{R}_+^{n+1}} e^{\lambda\omega(\xi)} | \int_{\mathbb{R}_+^{n+1}} e^{-i\langle \xi', y' \rangle} J_\beta(\xi_{n+1} y_{n+1}) (-\|y\|^2)^\alpha \phi(y) d\mu_\beta(y) | . \end{aligned}$$

Now, we can write $g(y) = (-\|y\|^2)^\alpha \phi(y)$ therefore above expression yields

$$\begin{aligned} p_{\alpha,\lambda}[\mathcal{F}_w\phi(\xi)] &= \sup_{\xi \in \mathbb{R}_+^{n+1}} e^{\lambda\omega(\xi)} | \int_{\mathbb{R}_+^{n+1}} e^{-i\langle \xi', y' \rangle} J_\beta(\xi_{n+1} y_{n+1}) g(y) d\mu_\beta(y) | \\ &\leq \sup_{\xi \in \mathbb{R}_+^{n+1}} e^{\lambda|\omega(\xi)|} | \int_{\mathbb{R}_+^{n+1}} e^{-i\langle \xi', y' \rangle} J_\beta(\xi_{n+1} y_{n+1}) g(y) d\mu_\beta(y) | . \end{aligned}$$

From [46] for every $\epsilon > 0$ there exist $C(\epsilon)$ such that

$$|\omega(\xi)| \leq \epsilon \|\xi\|^2 + C(\epsilon).$$

Therefore, we have

$$\begin{aligned} & p_{\alpha,\lambda}[\mathcal{F}_w\phi(\xi)] \\ & \leq \sup_{\xi \in \mathbb{R}_+^{n+1}} e^{\lambda(\epsilon\|\xi\|^2 + C(\epsilon))} \left| \int_{\mathbb{R}_+^{n+1}} e^{-i\langle \xi', y' \rangle} J_\beta(\xi_{n+1} y_{n+1}) g(y) d\mu_\beta(y) \right| \\ & \leq e^{\lambda C(\epsilon)} \sup_{\xi \in \mathbb{R}_+^{n+1}} e^{\lambda\epsilon\|\xi\|^2} \left| \int_{\mathbb{R}_+^{n+1}} e^{-i\langle \xi', y' \rangle} J_\beta(\xi_{n+1} y_{n+1}) g(y) d\mu_\beta(y) \right| \\ & \leq e^{\lambda C(\epsilon)} \sup_{\xi \in \mathbb{R}_+^{n+1}} \sum_{m=0}^{\infty} \frac{(\lambda\epsilon)^m}{m!} \|\xi\|^{2m} \left| \int_{\mathbb{R}_+^{n+1}} e^{-i\langle \xi', y' \rangle} J_\beta(\xi_{n+1} y_{n+1}) g(y) d\mu_\beta(y) \right| \\ & = e^{\lambda C(\epsilon)} \sum_{m=0}^{\infty} \frac{(\lambda\epsilon)^m}{m!} \sup_{\xi \in \mathbb{R}_+^{n+1}} \left| \int_{\mathbb{R}_+^{n+1}} \|\xi\|^{2m} e^{-i\langle \xi', y' \rangle} J_\beta(\xi_{n+1} y_{n+1}) g(y) d\mu_\beta(y) \right| \\ & = e^{\lambda C(\epsilon)} \sum_{m=0}^{\infty} \frac{(\lambda\epsilon)^m}{m!} \sup_{\xi \in \mathbb{R}_+^{n+1}} \left| \int_{\mathbb{R}_+^{n+1}} (\Delta_{W,\beta}^n)^m (e^{-i\langle \xi', y' \rangle} J_\beta(\xi_{n+1} y_{n+1})) g(y) d\mu_\beta(y) \right|. \end{aligned}$$

By integrating by parts, we get

$$\begin{aligned} & p_{\alpha,\lambda}[\mathcal{F}_w\phi(\xi)] \\ & \leq e^{\lambda C(\epsilon)} \sum_{m=0}^{\infty} \frac{(\lambda\epsilon)^m}{m!} \sup_{\xi \in \mathbb{R}_+^{n+1}} \left| (-1)^m \int_{\mathbb{R}_+^{n+1}} e^{-i\langle \xi', y' \rangle} J_\beta(\xi_{n+1} y_{n+1}) (\Delta_{W,\beta}^n)^m g(y) d\mu_\beta(y) \right| \\ & \leq e^{\lambda C(\epsilon)} \sum_{m=0}^{\infty} \frac{(\lambda\epsilon)^m}{m!} \sup_{\xi \in \mathbb{R}_+^{n+1}} \left| \int_{\mathbb{R}_+^{n+1}} e^{-i\langle \xi', y' \rangle} J_\beta(\xi_{n+1} y_{n+1}) (\Delta_{W,\beta}^n)^m g(y) d\mu_\beta(y) \right| \end{aligned}$$

Therefore, we have

$$\begin{aligned}
 & p_{\alpha,\lambda}[\mathcal{F}_w\phi(\xi)] \\
 & \leq e^{\lambda C(\epsilon)} \sum_{m=0}^{\infty} \frac{(\lambda\epsilon)^m}{m!} \left(\sup_{\xi \in \mathbb{R}_+^{n+1}} |e^{-i\langle \xi', y' \rangle} J_\beta(\xi_{n+1} y_{n+1})| \right) \left| \int_{\mathbb{R}_+^{n+1}} (\Delta_{W,\beta}^n)_y^m g(y) d\mu_\beta(y) \right| \\
 & \leq e^{\lambda C(\epsilon)} \sum_{m=0}^{\infty} \frac{(\lambda\epsilon)^m}{m!} \left| \int_{\mathbb{R}_+^{n+1}} e^{\lambda\omega(y)} (\Delta_{W,\beta}^n)_y^m g(y) e^{-\lambda\omega(y)} d\mu_\beta(y) \right| \\
 & \leq e^{\lambda C(\epsilon)} \sum_{m=0}^{\infty} \frac{(\lambda\epsilon)^m}{m!} \sup_{y \in \mathbb{R}_+^{n+1}} |e^{\lambda\omega(y)} (\Delta_{W,\beta}^n)_y^m g(y)| \left| \int_{\mathbb{R}_+^{n+1}} e^{-\lambda\omega(y)} d\mu_\beta(y) \right| \\
 & = e^{\lambda C(\epsilon)} \sum_{m=0}^{\infty} \frac{(\lambda\epsilon)^m}{m!} \left(\sup_{y \in \mathbb{R}_+^{n+1}} |(\Delta_{W,\beta}^n)_y^m g(y)| \right) \left| \int_{\mathbb{R}_+^{n+1}} e^{-\lambda\omega(y)} d\mu_\beta(y) \right| \\
 & = e^{\lambda C(\epsilon)} \sum_{m=0}^{\infty} \frac{(\lambda\epsilon)^m}{m!} p_{m,\lambda}(g(y)) \left| \int_{\mathbb{R}_+^{n+1}} e^{-\lambda\omega(y)} d\mu_\beta(y) \right| \\
 & < \infty.
 \end{aligned}$$

This shows that $\mathcal{F}_w(\phi) \in S_\omega(\mathbb{R}_+^{n+1})$.

Since by the inversion formula for the Weinstein transform we have $\mathcal{F}_w^{-1}\phi(y) = \mathcal{F}_w\phi(-y)$. Therefore \mathcal{F}_w is one-one. So that \mathcal{F}_w is an automorphism.

Definition 4.2.2. The generalized Sobolev type space $G_\omega^{p,s}$ is the set of all ultradistributions $u \in S'_\omega(\mathbb{R}_+^{n+1})$, such that its the Weinstein transform $(\mathcal{F}_w u)$ corresponds to a locally integrable function u over \mathbb{R}_+^{n+1} for which

$$\|u\|_{G_\omega^{p,s}} = \|e^{s\omega(\xi)}(\mathcal{F}_w u)(\xi)\|_p < \infty. \quad (4.2.1)$$

Theorem 4.2.3. The space $G_\omega^{p,s}$, $1 \leq p < \infty$ is complete.

Proof: Let $\{u_n\}$ be a Cauchy sequence in $G_\omega^{p,s}$. Then $e^{s\omega(\xi)}(\mathcal{F}_w u_n(\xi))$ is a Cauchy sequence in $L^p(\mathbb{R}_+^{n+1})$. Since $L^p(\mathbb{R}_+^{n+1})$ is complete, therefore there exists a function

$f \in L^p(\mathbb{R}_+^{n+1})$ such that

$$e^{s\omega(\xi)}(\mathcal{F}_w u_n)(\xi) \rightarrow f(\xi) \in L^p(\mathbb{R}_+^{n+1}) \text{ as } n \rightarrow \infty. \quad (4.2.2)$$

Let

$$g(\xi) = e^{-s\omega(\xi)} f(\xi), \quad s > 0. \quad (4.2.3)$$

Then $f \in S'_\omega(\mathbb{R}_+^{n+1})$ and $g(\xi) = e^{-s\omega(\xi)} f(\xi)$ are in $S'_\omega(\mathbb{R}_+^{n+1})$. Therefore, the inverse Weinstein transform of the above function exists. Using (4.2.3) we have

$$e^{s\omega(\xi)} g(\xi) = e^{s\omega(\xi)} \mathcal{F}_w(\mathcal{F}_w^{-1} g)(\xi) \in L^p(\mathbb{R}_+^{n+1}).$$

In view of (4.2.1), we get

$$\mathcal{F}_w^{-1} g \in G_\omega^{p,s}.$$

Let $\mathcal{F}_w^{-1} g = u$. Then

$$e^{s\omega(\xi)}(\mathcal{F}_w u)(\xi) = e^{s\omega(\xi)}(\mathcal{F}_w(\mathcal{F}_w^{-1} g))(\xi) = e^{s\omega(\xi)} g(\xi) = e^{s\omega(\xi)} e^{-s\omega(\xi)} f(\xi) = f(\xi).$$

This implies

$$e^{s\omega(\xi)}(\mathcal{F}_w u)(\xi) = f(\xi). \quad (4.2.4)$$

Using (4.2.2) and (4.2.4) we get $u_n \rightarrow u$ in $G_\omega^{p,s}$.

This shows that $G_\omega^{p,s}$ is complete.

Theorem 4.2.4. For $t > r$, $G_\omega^{p,t} \subseteq G_\omega^{p,r}$, the inclusion map is continuous.

Proof: Let $u \in G_\omega^{p,t}$ and $t > r$ then we take $t = r + \epsilon$, for some $\epsilon > 0$. Then

$$\begin{aligned} \|u\|_{G_\omega^{p,t}} &= \|e^{t\omega(\xi)}(\mathcal{F}_w u)(\xi)\|_p = \left(\int_{\mathbb{R}_+^{n+1}} |e^{t\omega(\xi)}(\mathcal{F}_w u)(\xi)|^p d\mu_\beta(\xi) \right)^{\frac{1}{p}} \\ &= \left(\int_{\mathbb{R}_+^{n+1}} |e^{(r+\epsilon)\omega(\xi)}(\mathcal{F}_w u)(\xi)|^p d\mu_\beta(\xi) \right)^{\frac{1}{p}} \\ &\geq \left(\int_{\mathbb{R}_+^{n+1}} |e^{r\omega(\xi)}(\mathcal{F}_w u)(\xi)|^p d\mu_\beta(\xi) \right)^{\frac{1}{p}} = \|e^{r\omega(\xi)}(\mathcal{F}_w u)(\xi)\|_p = \|u\|_{G_\omega^{p,r}}. \end{aligned}$$

Therefore,

$$\|u\|_{G_\omega^{p,r}} \leq \|u\|_{G_\omega^{p,t}}.$$

Hence

$$G_\omega^{p,t} \subseteq G_\omega^{p,r}.$$

Now let $\{u_n\} \rightarrow 0$ as $n \rightarrow \infty$ in $G_\omega^{p,t}$ then

$$\|u_n\|_{G_\omega^{p,t}} = \|e^{t\omega(\xi)}(\mathcal{F}_w u_n)(\xi)\|_p \rightarrow 0.$$

This implies that

$$\|u_n\|_{G_\omega^{p,r}} = \|e^{r\omega(\xi)}(\mathcal{F}_w u_n)(\xi)\|_p \rightarrow 0 \text{ as } n \rightarrow \infty.$$

This shows that $\{u_n\} \rightarrow 0$ in $G_\omega^{p,r}$ as $n \rightarrow \infty$.

This proves the continuity of the inclusion.

Theorem 4.2.5. *The space $S_\omega(\mathbb{R}_+^{n+1})$ is dense in $G_\omega^{p,s}$.*

Proof: Let $u \in G_\omega^{p,s}$. Then $e^{s\omega(\xi)}(\mathcal{F}_w u)(\xi) \in L^p(\mathbb{R}_+^{n+1})$. Since space $D_\omega(\mathbb{R}_+^{n+1})$ is dense in $L^p(\mathbb{R}_+^{n+1})$, then there exists a sequence of functions $\{\phi_n\}$ belonging to

$D_\omega(\mathbb{R}_+^{n+1})$ such that

$$\phi_n(\xi) \rightarrow e^{s\omega(\xi)}(F_w u)(\xi) \text{ in } L^p(\mathbb{R}_+^{n+1}). \quad (4.2.5)$$

Now, define

$$\psi_n(x) = \mathcal{F}_w^{-1}(e^{-s\omega(\xi)}\phi_n)(x). \quad (4.2.6)$$

Since $\phi_n \in D_\omega(\mathbb{R}_+^{n+1})$, therefore $(e^{-s\omega(\xi)}\phi_n) \in D_\omega(\mathbb{R}_+^{n+1})$, as $D_\omega(\mathbb{R}_+^{n+1})$ is subspace of $S_\omega(\mathbb{R}_+^{n+1})$, $(e^{-s\omega(\xi)}\phi_n)$ also in $S_\omega(\mathbb{R}_+^{n+1})$. Hence, its inverse Weinstein transform exists, so that

$$\psi_n(x) = \mathcal{F}_w^{-1}(e^{-s\omega(\xi)}\phi_n)(x) \in H_\omega(\mathbb{R}_+^{n+1}).$$

From (4.2.1) we can get

$$\begin{aligned} \|\psi_n - u\|_{G_\omega^{p,s}} &= \|e^{s\omega(\xi)}(\mathcal{F}_w(\psi_n - u)(\xi))\|_p \\ &= \|e^{s\omega(\xi)}(\mathcal{F}_w\psi_n)(\xi) - e^{s\omega(\xi)}(\mathcal{F}_w u)(\xi)\|_p. \end{aligned}$$

In view of (4.2.6), we have

$$\|\psi_n - u\|_{G_\omega^{p,s}} = \|\phi_n - e^{s\omega(\xi)}(\mathcal{F}_w u)(\xi)\|_p.$$

Using (4.2.5), we find

$$\|\psi_n - u\|_{G_\omega^{p,s}} = \|\phi_n - e^{s\omega(\xi)}(\mathcal{F}_w u)(\xi)\|_p \rightarrow 0 \text{ as } n \rightarrow \infty.$$

Hence,

$$\psi_n \rightarrow u \text{ in } G_\omega^{p,s}.$$

This shows that the space $S_\omega(\mathbb{R}_+^{n+1})$ is dense in $G_\omega^{p,s}$.

Lemma 4.2.6. *Let s_1, s and s_2 be three real numbers with $s_1 < s < s_2$ then for every $\epsilon > 0$ there exists $C(\epsilon) > 0$ such that*

$$e^{s\omega(\xi)} \leq \epsilon e^{s_2\omega(\xi)} + C(\epsilon)e^{s_1\omega(\xi)}. \quad (4.2.7)$$

Proof: The proof of above Lemma can be obtained from [45].

Theorem 4.2.7. *Let s_1, s and s_2 be three real numbers with $s_1 < s < s_2$ then for every $\epsilon > 0$ there exist $C(\epsilon) > 0$ such that*

$$\|u\|_{G_\omega^{p,s}} \leq \epsilon \|u\|_{G_\omega^{p,s_2}} \leq C(\epsilon) \|u\|_{G_\omega^{p,s_1}} \quad \forall u \in \|u\|_{G_\omega^{p,s_2}}$$

Proof: Let $u \in G_\omega^{p,s_2}$. Then by Theorem (4.2.4), we have

$$u \in G_\omega^{p,s} \text{ and } u \in G_\omega^{p,s_1}.$$

Then

$$\begin{aligned} \|u\|_{u \in G_\omega^{p,s}} &= \|e^{s\omega(\xi)}(\mathcal{F}_w u)(\xi)\|_p \\ &= \left(\int_{\mathbb{R}_+^{n+1}} |e^{s\omega(\xi)}(\mathcal{F}_w u)(\xi)|^p d\mu_\beta(\xi) \right)^{\frac{1}{p}}. \end{aligned}$$

By using (4.2.7), we get

$$\|u\|_{u \in G_\omega^{p,s}} \leq \left(\int_{\mathbb{R}_+^{n+1}} |(\epsilon e^{s_2\omega(\xi)} + C(\epsilon)e^{s_1\omega(\xi)})(\mathcal{F}_w u)(\xi)|^p d\mu_\beta(\xi) \right)^{\frac{1}{p}}$$

So that we find that

$$\begin{aligned}
 \|u\|_{u \in G_\omega^{p,s}} &\leq \left(\int_{\mathbb{R}_+^{n+1}} |\epsilon e^{s_2 \omega(\xi)} (\mathcal{F}_w u)(\xi)|^p d\mu_\beta(\xi) \right)^{\frac{1}{p}} \\
 &\quad + \left(\int_{\mathbb{R}_+^{n+1}} |C(\epsilon) e^{s_1 \omega(\xi)} \mathcal{F}_w u(\xi)|^p d\mu_\beta(\xi) \right)^{\frac{1}{p}} \\
 &\leq \epsilon \left(\int_{\mathbb{R}_+^{n+1}} |e^{s_2 \omega(\xi)} (\mathcal{F}_w u)(\xi)|^p d\mu_\beta(\xi) \right)^{\frac{1}{p}} \\
 &\quad + C(\epsilon) \left(\int_{\mathbb{R}_+^{n+1}} |e^{s_1 \omega(\xi)} \mathcal{F}_w u(\xi)|^p d\mu_\beta(\xi) \right)^{\frac{1}{p}} \\
 &= \epsilon \|e^{s_2 \omega(\xi)} (\mathcal{F}_w u)(\xi)\|_p + C(\epsilon) \|e^{s_1 \omega(\xi)} (\mathcal{F}_w u)(\xi)\|_p \\
 &= \epsilon \|u\|_{u \in G_\omega^{p,s_2}} + C(\epsilon) \|u\|_{u \in G_\omega^{p,s_1}}.
 \end{aligned}$$

This shows that

$$\|u\|_{G_\omega^{p,s}} \leq \epsilon \|u\|_{G_\omega^{p,s_2}} \leq C(\epsilon) \|u\|_{G_\omega^{p,s_1}}.$$

This proves the theorem.

Theorem 4.2.8. *Let $s < 0$, then $L^2(\mathbb{R}_+^{n+1}) \subset G_\omega^{2,s}$.*

Proof: Let $\phi \in L^2(\mathbb{R}_+^{n+1})$. Then from (1.4.3), we find

$$\int_{\mathbb{R}_+^{n+1}} |(\mathcal{F}_w \phi)(\xi)|^2 d\mu_\beta(\xi) = \int_{\mathbb{R}_+^{n+1}} |\phi(x)|^2 d\mu_\beta(x).$$

Therefore, above implies

$$\left(\int_{\mathbb{R}_+^{n+1}} |(\mathcal{F}_w \phi)(\xi)|^2 d\mu_\beta(\xi) \right)^{\frac{1}{2}} = \left(\int_{\mathbb{R}_+^{n+1}} |\phi(x)|^2 d\mu_\beta(x) \right)^{\frac{1}{2}}.$$

Therefore, we have

$$\|\mathcal{F}_w \phi\|_{L^2(\mathbb{R}_+^{n+1})} = \|\phi\|_{L^2(\mathbb{R}_+^{n+1})}. \quad (4.2.8)$$

Hence for $s < 0$, we have

$$\begin{aligned} \|\phi\|_{G_\omega^{2,s}} &= \|e^{s\omega(\xi)}(\mathcal{F}_w\phi)(\xi)\|_{L^2(\mathbb{R}_+^{n+1})} = \left(\int_{\mathbb{R}_+^{n+1}} |e^{s\omega(\xi)}(\mathcal{F}_w\phi)(\xi)|^2 d\mu_\beta(\xi) \right)^{\frac{1}{2}} \\ &\leq \left(\int_{\mathbb{R}_+^{n+1}} |(\mathcal{F}_w\phi)(\xi)|^2 d\mu_\beta(\xi) \right)^{\frac{1}{2}} = \|(\mathcal{F}_w\phi)(\xi)\|_{L^2(\mathbb{R}_+^{n+1})}. \end{aligned}$$

By using (4.2.8), we get

$$\|\phi\|_{G_\omega^{2,s}} \leq \|(\mathcal{F}_w\phi)(\xi)\|_{L^2(\mathbb{R}_+^{n+1})}.$$

Therefore

$$\|\phi\|_{G_\omega^{2,s}} < \infty.$$

Hence, above implies

$$\phi \in G_\omega^{2,s}.$$

This shows that $L^2(\mathbb{R}_+^{n+1}) \subset G_\omega^{2,s}$.

4.3 The generalized Hankel potential associated with the Weinstein transform

In this section, we discussed various properties of the Hankel potential associated with the Weinstein transform.

Let $a(x) \neq 0$ be a multiplier in $S_\omega(\mathbb{R}_+^{n+1})$ such that

$$\mathcal{F}_w^{-1}(a^k)(x) \in L^1(\mathbb{R}_+^{n+1}) \quad \forall \quad k = 0, 1, 2, 3, \dots \quad (4.3.1)$$

Then, for $f \in S'_\omega(\mathbb{R}_+^{n+1})$, the products fa^k and $fa^{-k} \in S'_\omega(\mathbb{R}_+^{n+1})$ are defined by the following relations.

$$\langle f a^k, \phi \rangle = \langle f, a^k \phi \rangle \quad \forall \phi \in S_\omega(\mathbb{R}_+^{n+1}) \quad (4.3.2)$$

and

$$\langle f a^{-k}, a^k \phi \rangle = \langle f, \phi \rangle \quad \forall \phi \in S_\omega(\mathbb{R}_+^{n+1}). \quad (4.3.3)$$

Now, the generalized Hankel potential associated with the Weinstein transform is

$$\mathcal{H}^k u = \mathcal{F}_w^{-1}(a^{-k}(\mathcal{F}_w u)) \quad \forall u \in S'_\omega(\mathbb{R}_+^{n+1}). \quad (4.3.4)$$

Since the generalized Weinstein transformation \mathcal{F}'_w corresponds to a continuous linear mapping between $S'_\omega(\mathbb{R}_+^{n+1})$ onto $S'_\omega(\mathbb{R}_+^{n+1})$ and inverse of the Weinstein transform holds for $(\mathcal{F}'_w)^{-1}$ as well . Therefore we conclude that the generalized Hankel potential associated with the Weinstein transform (4.3.5) is a continuous linear mapping of $S'_\omega(\mathbb{R}_+^{n+1})$ onto $S'_\omega(\mathbb{R}_+^{n+1})$. From (1.4.1) and (1.4.6) we have

$$(\mathcal{H}^k u)(x) = \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1} \xi_{n+1}) a^{-k}(\xi) (\mathcal{F}_w u)(\xi) d\mu_\beta(\xi). \quad (4.3.5)$$

From [84, p.10], the generalized Hankel potential associated with the Weinstein transform is a pseudo differential operator with symbol $a^{-k}(\xi)$

Theorem 4.3.1. *Let $u \in S'_\omega(\mathbb{R}_+^{n+1})$. Then for any $m, l \in \mathbb{Z}$ we have*

(i) $\mathcal{H}^l \mathcal{H}^m u = \mathcal{H}^{l+m} u.$

(ii) $\mathcal{H}^0 u = u.$

Proof:(i): Let $u \in S'_\omega(\mathbb{R}_+^{n+1})$ then from (4.3.4), we have

$$(\mathcal{H}^l \mathcal{H}^m u)(x) = \mathcal{F}_w^{-1} \left(a^{-l} (\mathcal{F}_w (\mathcal{H}^m u)) \right) (x).$$

Using (1.4.1), we get

$$\begin{aligned} (\mathcal{H}^l \mathcal{H}^m u)(x) &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1} \xi_{n+1}) (a^{-l} \mathcal{F}_w (\mathcal{H}^m u))(\xi) d\mu_\beta(\xi) \\ &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1} \xi_{n+1}) a^{-l}(\xi) \mathcal{F}_w (\mathcal{H}^m u)(\xi) d\mu_\beta(\xi). \end{aligned}$$

In view of (4.3.4), we have

$$(\mathcal{H}^m u)(\xi) = (\mathcal{F}_w^{-1} (a^{-m} (\mathcal{F}_w u)))(\xi).$$

Then, by the Weinstein transform we find

$$\mathcal{F}_w (\mathcal{H}^m u)(\xi) = a^{-m}(\xi) ((\mathcal{F}_w u))(\xi). \quad (4.3.6)$$

Using (4.3.6), we find

$$\begin{aligned} (\mathcal{H}^l \mathcal{H}^m u)(x) &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1} \xi_{n+1}) a^{-l}(\xi) a^{-m}(\xi) (\mathcal{F}_w u)(\xi) d\mu_\beta(\xi) \\ &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1} \xi_{n+1}) a^{-(l+m)}(\xi) (\mathcal{F}_w u)(\xi) d\mu_\beta(\xi) \\ &= \mathcal{F}_w^{-1} \left(a^{-(l+m)}(\xi) \mathcal{F}_w u(\xi) \right) (x) \\ &= (\mathcal{H}^{l+m} u)(x). \end{aligned}$$

Therefore,

$$\mathcal{H}^l \mathcal{H}^m u = \mathcal{H}^{l+m} u.$$

(ii): Now, we have

$$\begin{aligned} (\mathcal{H}^0 u)(x) &= \mathcal{F}_w^{-1}(a^0(\mathcal{F}_w u))(x) \\ &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1} \xi_{n+1})(a^{-0}(\xi))(\mathcal{F}_w(u))(\xi) d\mu_\beta(\xi) \\ &= \mathcal{F}_w^{-1}(\mathcal{F}_w u)(x) \\ &= u(x). \end{aligned}$$

This shows that

$$\mathcal{H}^0 u = u.$$

The Space $W_\omega^{m,p}(\mathbb{R}_+^{n+1})$

In this section we introduce the L^p -space of all Hankel potential \mathcal{H}^k which is represented by the $W_\omega^{m,p}(\mathbb{R}_+^{n+1})$ and exploiting the theory of Weinstein transform, many properties of this space will be discussed.

Definition 4.3.2. Let $m \in \mathbb{Z}$ and $1 < p < \infty$, the space $W_\omega^{m,p}(\mathbb{R}_+^{n+1})$ is the set of all those ultradistributions $u \in S'_\omega(\mathbb{R}_+^{n+1})$ for which

$$\mathcal{H}^{-m} u \in L^p(\mathbb{R}_+^{n+1}). \quad (4.3.7)$$

The norm of this space is defined by

$$\|u\|_{m,p,\omega} = \|\mathcal{H}^{-m} u\|_p = \left(\int_{\mathbb{R}_+^{n+1}} |(\mathcal{H}^{-m} u)(\xi)|^p d\mu_\beta(\xi) \right)^{\frac{1}{p}}. \quad (4.3.8)$$

In the next theorem we are going to prove that the space $W_\omega^{m,p}(\mathbb{R}_+^{n+1})$ is a Banach space.

Theorem 4.3.3. *The space $W_\omega^{m,p}(\mathbb{R}_+^{n+1})$ is a Banach space with respect to the norm $\|\cdot\|_{m,p,\omega}$.*

Proof: Let $\{u_k\}$ be a Cauchy sequence in $W_\omega^{m,p}(\mathbb{R}_+^{n+1})$. Then $\{\mathcal{H}^{-m}u_k\}$ is a Cauchy sequence in $L^p(\mathbb{R}_+^{n+1})$. Since $L^p(\mathbb{R}_+^{n+1})$ is complete, therefore there exists a function $u \in L^p(\mathbb{R}_+^{n+1})$ such that

$$\mathcal{H}^{-m}u_k \rightarrow u \in L^p(\mathbb{R}_+^{n+1}) \text{ as } k \rightarrow \infty. \quad (4.3.9)$$

Take $v = \mathcal{H}^m u$. Then from (4.3.5), we have

$$\begin{aligned} v(x) &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1} \xi_{n+1}) a^{-m}(\xi) (\mathcal{F}_w u)(\xi) d\mu_\beta(\xi) \\ &= \mathcal{F}_w^{-1}(a^{-m}(\mathcal{F}_w u))(x). \end{aligned}$$

Therefore, by the property of the Weinstein transform (1.4.6), we get

$$u(x) = \mathcal{F}_w^{-1}(a^m(\mathcal{F}_w v))(x) = (\mathcal{H}^{-m}v)(x) \in L^p(\mathbb{R}_+^{n+1}). \quad (4.3.10)$$

Hence, we find

$$v \in W_p^{m,p}.$$

From (4.3.9) and (4.3.10), it follows that

$$u_k \rightarrow v \in W_p^{m,p} \text{ as } k \rightarrow \infty.$$

This proves that the space $W_\omega^{m,p}(\mathbb{R}_+^{n+1})$ is a Banach space.

Theorem 4.3.4. *The generalized Hankel potential \mathcal{H}^t is an isometry of $W_\omega^{m,p}(\mathbb{R}_+^{n+1})$ onto $W_\omega^{m+t,p}(\mathbb{R}_+^{n+1})$; and we have*

$$\|\mathcal{H}^t \phi\|_{m+t,p,\omega} = \|\phi\|_{m,p,\omega}, \quad \forall \phi \in W_\omega^{m,p}(\mathbb{R}_+^{n+1}). \quad (4.3.11)$$

Proof: Let $\phi \in W_\omega^{m,p}$, then from (4.3.8) we have

$$\|\mathcal{H}^t \phi\|_{m+t,p,\omega} = \|\mathcal{H}^{-(m+t)}(\mathcal{H}^t \phi)\|_p. \quad (4.3.12)$$

Now, we have

$$(\mathcal{H}^{-(m+t)}(\mathcal{H}^t \phi))(x) = \mathcal{F}_w^{-1}(a^{m+t}(x)\mathcal{F}_w(\mathcal{H}^t \phi))(x)$$

By inversion formula of the Weinstein transform (1.4.4), above expression yields

$$(\mathcal{H}^{-(m+t)}(\mathcal{H}^t \phi))(x) = \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1}\xi_{n+1}) a^{m+t}(\xi) \mathcal{F}_w(\mathcal{H}^t \phi)(\xi) d\mu_\beta(\xi).$$

Taking (4.3.6), we get

$$\begin{aligned} & (\mathcal{H}^{-(m+t)}(\mathcal{H}^t \phi))(x) \\ &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1}\xi_{n+1}) a^{m+t}(\xi) a^{-t}(\xi) (\mathcal{F}_w \phi)(\xi) d\mu_\beta(\xi) \\ &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1}\xi_{n+1}) a^m(\xi) (\mathcal{F}_w \phi)(\xi) d\mu_\beta(\xi) \\ &= (\mathcal{H}^{-m} \phi)(x). \end{aligned}$$

Therefore, from above expression we find that

$$(\mathcal{H}^{-(m+t)}(\mathcal{H}^t \phi)) = (\mathcal{H}^{-m} \phi).$$

In view of (4.3.12), we have

$$\begin{aligned} \|\mathcal{H}^t \phi\|_{m+t,p,\omega} &= \|(\mathcal{H}^{-m} \phi)\|_p \\ &= \|\phi\|_{m,p,\omega}. \end{aligned}$$

Let $\phi \in W_\omega^{m+t,p}(\mathbb{R}_+^{n+1})$. Then from (4.3.7), we obtain

$$\mathcal{H}^{-(m+t)} \phi \in L^p(\mathbb{R}_+^{n+1}).$$

Again, we get

$$\begin{aligned} (\mathcal{H}^{-(m+t)} \phi)(x) &= \mathcal{F}_w^{-1}(a^{m+t}(\xi)(\mathcal{F}_w \phi)(\xi))(x) \\ &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1} \xi_{n+1}) a^{m+t}(\xi) (\mathcal{F}_w \phi)(\xi) d\mu_\beta(\xi) \\ &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1} \xi_{n+1}) a^m(\xi) a^t(\xi) (\mathcal{F}_w \phi)(\xi) d\mu_\beta(\xi). \end{aligned}$$

Then applying (4.3.6), we find

$$\begin{aligned} (\mathcal{H}^{-(m+t)} \phi)(x) &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1} \xi_{n+1}) a^m(\xi) \mathcal{F}_w(\mathcal{H}^{-t} \phi)(\xi) d\mu_\beta(\xi) \\ &= (\mathcal{H}^{-m}(\mathcal{H}^{-t} \phi))(x). \end{aligned}$$

This show that

$$\mathcal{H}^{-(m+t)} \phi = \mathcal{H}^{-m}(\mathcal{H}^{-t} \phi) \in L^p(\mathbb{R}_+^{n+1}).$$

Hence,

$$(\mathcal{H}^{-t} \phi) \in W_\omega^{m,p}.$$

And, we have to obtain

$$\begin{aligned}
 & \mathcal{H}^t(\mathcal{H}^{-t}\phi)(x) \\
 &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1}\xi_{n+1}) a^{-t}(\xi) \mathcal{F}_w(\mathcal{H}^{-t}\phi)(\xi) d\mu_\beta(\xi) \\
 &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1}\xi_{n+1}) a^{-t}(\xi) b^{-t}(\xi) \mathcal{F}_w\phi(\xi) d\mu_\beta(\xi) \\
 &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1}\xi_{n+1}) (\mathcal{F}_w\phi)(\xi) d\mu_\beta(\xi) \\
 &= \mathcal{F}^{-1}((\mathcal{F}_w\phi)(\xi))(x) = \phi(x).
 \end{aligned}$$

Therefore, for each $\phi \in W_\omega^{m,p}$, there exists an $\mathcal{H}^{-t}\phi \in W_\omega^{m+t,p}$ such that

$$\mathcal{H}^t\mathcal{H}^{-t}\phi = \phi.$$

Hence, \mathcal{H}^t is onto.

Theorem 4.3.5. *Let $1 < p < \infty$ and $m = 0, 1, 2, \dots$ then,*

$$\|\mathcal{H}^m\phi\| \leq C\|\phi\|_p \quad \forall \phi \in S_\omega(\mathbb{R}_+^{n+1}). \quad (4.3.13)$$

Proof: From (1.4.13), we have

$$\mathcal{F}_w(f\#_\beta\phi)(\xi) = (\mathcal{F}_wf)(\xi)(\mathcal{F}_w\phi)(\xi). \quad (4.3.14)$$

Take

$$a^{-m}(\xi) = \mathcal{F}_w(f(t))(\xi). \quad (4.3.15)$$

Then,

$$\begin{aligned} f(t) &= \mathcal{F}_w^{-1}(a^{-m}(\xi))(t) = \mathcal{F}_w(a^{-m}(\xi))(-t) \\ &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle t', \xi' \rangle} J_\beta(t_{n+1}\xi_{n+1}) a^{-m}(\xi) d\mu_\beta(\xi) \\ &= \mathcal{F}_w^{-1}(a^{-m})(t). \end{aligned}$$

Utilising (4.3.1), it follows that

$$f(t) = \mathcal{F}_w^{-1}(a^{-m})(t) \in L^1(\mathbb{R}_+^{n+1}).$$

In view of (4.3.14), we have

$$\begin{aligned} \mathcal{F}_w(f\#_\beta\phi)(\xi) &= \mathcal{F}_w(\mathcal{F}_w^{-1}(a^{-m}(t))(\xi)(\mathcal{F}_w\phi)(\xi)) \\ &= a^{-m}(\xi)(\mathcal{F}_w\phi)(\xi). \end{aligned}$$

Therefore,

$$\begin{aligned} (f\#_\beta\phi)(t) &= \mathcal{F}_w^{-1}(a^{-m}(\xi)(\mathcal{F}_w\phi)(\xi))(t) \\ &= \mathcal{F}_w^{-1}(a^{-m}\mathcal{F}_w\phi)(t) \\ &= \mathcal{H}^m\phi. \end{aligned}$$

Then

$$\|\mathcal{H}^m\phi\|_p = \|f\#_\beta\phi\|_p.$$

Using (1.4.14), we obtain

$$\|\mathcal{H}^m\phi\|_p \leq \|f\|_1 \|\phi\|_p \leq C\|\phi\|_p.$$

This proves the theorem.

Theorem 4.3.6. *Let $1 < p < \infty$ and $m \leq l$. Then*

$$W_{\omega}^{l,p}(\mathbb{R}_{+}^{n+1}) \subseteq W_{\omega}^{m,p}(\mathbb{R}_{+}^{n+1}),$$

and

$$\|\phi\|_{m,p,\omega} \leq C\|\phi\|_{l,p,\omega}.$$

Proof: Let $\phi \in W_{\omega}^{l,p}(\mathbb{R}_{+}^{n+1})$. Then from (4.3.7), we have

$$\mathcal{H}^{-l}\phi \in L^p(\mathbb{R}_{+}^{n+1}).$$

Then

$$\begin{aligned} (\mathcal{H}^{-m}\phi)(x) &= \mathcal{F}_w^{-1}(a^m \mathcal{F}_w \phi)(x) \\ &= \int_{\mathbb{R}_{+}^{n+1}} e^{i\langle x', \xi' \rangle} J_{\beta}(x_{n+1} \xi_{n+1}) a^m(\xi) (\mathcal{F}_w \phi)(\xi) d\mu_{\beta}(\xi) \\ &= \int_{\mathbb{R}_{+}^{n+1}} e^{i\langle x', \xi' \rangle} J_{\beta}(x_{n+1} \xi_{n+1}) a^{-(l-m)}(\xi) a^l(\xi) (\mathcal{F}_w \phi)(\xi) d\mu_{\beta}(\xi). \end{aligned}$$

From (4.3.6), we obtain

$$(\mathcal{H}^{-m}\phi)(x) = \int_{\mathbb{R}_{+}^{n+1}} e^{i\langle x', \xi' \rangle} J_{\beta}(x_{n+1} \xi_{n+1}) a^{-(l-m)}(\xi) \mathcal{F}_w(\mathcal{H}^l \phi)(\xi) d\mu_{\beta}(\xi).$$

From (4.3.4), we have

$$(\mathcal{H}^{-m}\phi)(x) = \mathcal{H}^{l-m}(\mathcal{H}^{-l}\phi)(x).$$

Therefore, in view of (4.3.7), we have

$$\|\phi\|_{m,p,\omega} = \|\mathcal{H}^{-m}\phi\|_p = \|\mathcal{H}^{l-m}(\mathcal{H}^{-l}\phi)\|_p.$$

Using (4.3.13), we get

$$\|\phi\|_{m,p,\omega} \leq C\|\mathcal{H}^{-l}\phi\|_p = \|\phi\|_{l,p,\omega}.$$

Hence,

$$\|\phi\|_{m,p,\omega} \leq \|\phi\|_{l,p,\omega}.$$

This proves the theorem.

Theorem 4.3.7. *Let $1 < p < \infty$ and for any symbol $a^m, m \in \mathbb{Z}$, define the pseudo differential operator*

$$(A(D)u)(x) = \mathcal{F}_w^{-1}(a^m \mathcal{F}_w u)(x).$$

Then

$$A(D) : W_\omega^{m,p}(\mathbb{R}_+^{n+1}) \longrightarrow W_\omega^{0,p}(\mathbb{R}_+^{n+1})$$

is a bounded linear operator.

Proof: Consider the following operators;

$$\mathcal{H}^{-l} : W_\omega^{l,p}(\mathbb{R}_+^{n+1}) \longrightarrow W_\omega^{0,p}(\mathbb{R}_+^{n+1}),$$

$$A(D)\mathcal{H}^m : W_\omega^{0,p}(\mathbb{R}_+^{n+1}) \longrightarrow W_\omega^{0,p}(\mathbb{R}_+^{n+1})$$

$$\mathcal{H}^{l-m} : W_\omega^{0,p}(\mathbb{R}_+^{n+1}) \longrightarrow W_\omega^{l-m,p}(\mathbb{R}_+^{n+1}).$$

We have to show that operator \mathcal{H}^{-l} is bounded. Now we take $\phi \in W_{\omega}^{l,p}$, then

$$\|\phi\|_{l,p,\omega} = \|\mathcal{H}^{-l}\phi\|_p < \infty, \quad (4.3.16)$$

and

$$\|\mathcal{H}^{-l}\phi\|_{0,p,\omega} = \|\mathcal{H}^0(\mathcal{H}^{-l}\phi)\|_p. \quad (4.3.17)$$

Now,

$$(\mathcal{H}^0(\mathcal{H}^{-l}\phi))(x) = \mathcal{F}_w^{-1}(a^0 \mathcal{F}_w(\mathcal{H}^{-l}\phi))(x).$$

By the inverse Weinstein transform (1.4.4), we get

$$\begin{aligned} (\mathcal{H}^0(\mathcal{H}^{-l}\phi))(x) &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_{\beta}(x_{n+1}\xi_{n+1}) a^0(\xi) \mathcal{F}_w(\mathcal{H}^{-l}\phi)(\xi) d\mu_{\beta}(\xi) \\ &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_{\beta}(x_{n+1}\xi_{n+1}) \mathcal{F}_w(\mathcal{H}^{-l}\phi)(\xi) d\mu_{\beta}(\xi) \\ &= \mathcal{F}_w^{-1}(\mathcal{F}_w(\mathcal{H}^{-l}\phi))(x) = (\mathcal{H}^{-l}\phi)(x). \end{aligned}$$

Therefore,

$$\mathcal{H}^0(\mathcal{H}^{-l}\phi) = (\mathcal{H}^{-l}\phi). \quad (4.3.18)$$

In veiw of (4.3.17), we get

$$\|\mathcal{H}^{-l}\phi\|_{0,p,\omega} = \|(\mathcal{H}^{-l}\phi)\|_p < \infty.$$

We have to show that the second operator is bounded. Let $\phi \in W_{\omega}^{0,p}$. Then

$$\|\phi\|_{0,p,\omega} = \|\mathcal{H}^0\phi\|_p = \|\phi\|_p < \infty, \quad (4.3.19)$$

and

$$\|A(D)\mathcal{H}^m\phi\|_{0,p,\omega} = \|A(D)\mathcal{H}^m\phi\|_p. \quad (4.3.20)$$

Now, we have to calculate

$$\begin{aligned} (A(D)\mathcal{H}^m\phi)(x) &= \mathcal{F}_w^{-1}(a^m(\xi)\mathcal{F}_w(\mathcal{H}^m\phi))(x) \\ &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1}\xi_{n+1}) a^m(\xi) \mathcal{F}_w(\mathcal{H}^m\phi)(\xi) d\mu_\beta(\xi). \end{aligned}$$

Therefore, using (4.3.6), we can obtain

$$\begin{aligned} (A(D)\mathcal{H}^m\phi)(x) &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1}\xi_{n+1}) a^m(\xi) a^{-m}(\xi) (\mathcal{F}_w\phi)(\xi) d\mu_\beta(\xi) \\ &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_\beta(x_{n+1}\xi_{n+1}) (\mathcal{F}_w\phi)(\xi) d\mu_\beta(\xi) \\ &= \mathcal{F}_w^{-1}(\mathcal{F}_w\phi)(x) = \phi(x). \end{aligned}$$

Using (4.3.19) and (4.3.20) we get

$$\|A(D)\mathcal{H}^m\phi\|_{0,p,\omega} = \|\phi\|_p < \infty.$$

Third operator \mathcal{H}^{l-m} , is bounded by the following way:

$$\|\mathcal{H}^{l-m}\phi\|_{l-m,p,\omega} = \|\mathcal{H}^{-(l-m)}\mathcal{H}^{l-m}\phi\|_p = \|\phi\|_p < \infty.$$

Hence, the product $\mathcal{H}^{l-m}A(D)\mathcal{H}^{m-l}$ is a bounded linear operator from $W_\omega^{l,p}(\mathbb{R}_+^{n+1})$ into $W_\omega^{l-m,p}(\mathbb{R}_+^{n+1})$.

Theorem 4.3.8. For $m \in \mathbb{Z}$ and symbol a^m , the pseudo differential operator

$$A(D) = \mathcal{F}_w^{-1}a^m\mathcal{F}_w : W_\omega^{l,p}(\mathbb{R}_+^{n+1}) \longrightarrow W_\omega^{l-m,p}(\mathbb{R}_+^{n+1})$$

is a bounded linear operator.

Proof: Let $\phi \in W_{\omega}^{l,p}(\mathbb{R}_+^{n+1})$. Then

$$\|\phi\|_{l,p,\omega} = \|\mathcal{H}^{-l}\phi\|_p < \infty,$$

and

$$\|A(D)\phi\|_{l-m,p,\omega} = \|\mathcal{H}^{m-l}(A(D)\phi)\|_p < \infty. \quad (4.3.21)$$

Therefore, we have

$$\begin{aligned} \mathcal{H}^{m-l}(A(D)\phi)(x) &= \mathcal{F}_w^{-1}(a^{-(m-l)}\mathcal{F}_w A(D)\phi)(x) \\ &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_{\beta}(x_{n+1}\xi_{n+1}) a^{-(m-l)}(\xi) \mathcal{F}_w(A(D)\phi)(\xi) d\mu_{\beta}(\xi) \\ &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_{\beta}(x_{n+1}\xi_{n+1}) a^{-(m-l)}(\xi) \mathcal{F}_w(\mathcal{F}_w^{-1} a^m \mathcal{F}_w \phi)(\xi) d\mu_{\beta}(\xi) \\ &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_{\beta}(x_{n+1}\xi_{n+1}) a^{-(m-l)}(\xi) a^m(\xi) (\mathcal{F}_w \phi)(\xi) d\mu_{\beta}(\xi) \\ &= \int_{\mathbb{R}_+^{n+1}} e^{i\langle x', \xi' \rangle} J_{\beta}(x_{n+1}\xi_{n+1}) a^l(\xi) (\mathcal{F}_w \phi)(\xi) d\mu_{\beta}(\xi) \end{aligned}$$

By (1.4.6), we have

$$\begin{aligned} \mathcal{H}^{m-l}(A(D)\phi)(x) &= \mathcal{F}_w^{-1}(a^l \mathcal{F}_w \phi)(x) \\ &= \mathcal{H}^{-l}\phi(x). \end{aligned}$$

Therefore, we have

$$\mathcal{H}^{m-l}(A(D)\phi)(x) = \mathcal{H}^{-l}\phi(x). \quad (4.3.22)$$

Using (4.3.21) and (4.3.22) the above expression yields

$$\|A(D)\phi\|_{l-m,p,\omega} = \|\mathcal{H}^{-l}\phi\|_p < \infty.$$

Above shows that the mapping

$$A(D) : W_{\omega}^{l,p}(\mathbb{R}_+^{n+1}) \longrightarrow W_{\omega}^{l-m,p}(\mathbb{R}_+^{n+1})$$

is a bounded linear operator.
